# Section A General Principles

# Units, Mathématics and Physical Quantities

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(Section 1.2.10)

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This reference section provides (a) a statement of the International System (SI) of Units, with conversion factors; (b) basic mathematical functions, series and tables; and (c) some physical properties of materials.

# 1.1 International unit system

The International System of Units (SI) is a metric system giving a fully coherent set of units for science, technology and engineering, involving no conversion factors. The starting point is the selection and definition of a minimum set of independent 'base' units. From these, 'derived' units are obtained by forming products or quotients in various combinations, again without numerical factors. For convenience, certain combinations are given shortened names. A single SI unit of energy (joule = kilogram metre-squared per second-squared) is, for example, applied to energy of any kind, whether it be kinetic, potential, electrical, thermal, chemical..., thus unifying usage throughout science and technology.

The SI system has seven *base* units, and two *supplement-ary* units of angle. Combinations of these are *derived* for all other units. Each physical quantity has a quantity symbol (e.g. *m* for mass, *P* for power) that represents it in physical equations, and a unit symbol (e.g. kg for kilogram, W for watt) to indicate its SI unit of measure.

# 1.1.1 Base units

Definitions of the seven base units have been laid down in the following terms. The quantity symbol is given in italic, the unit symbol (with its standard abbreviation) in roman type. As measurements become more precise, changes are occasionally made in the definitions.

*Length: l*, metre (m) The metre was defined in 1983 as the length of the path travelled by light in a vacuum during a time interval of 1/299792458 of a second.

*Mass: m*, kilogram (kg) The mass of the international prototype (a block of platinum preserved at the International Bureau of Weights and Measures, Sèvres).

*Time: t*, second (s) The duration of 9 192 631 770 periods of the radiation corresponding to the transition between the two hyperfine levels of the ground state of the caesium-133 atom.

*Electric current: i*, ampere (A) The current which, maintained in two straight parallel conductors of infinite length, of negligible circular cross-section and 1 m apart in vacuum, produces a force equal to  $2 \times 40^{-7}$  newton per metre of length.

Thermodynamic temperature: T, kelvin (K) The fraction 1/273.16 of the thermodynamic (absolute) temperature of the triple point of water.

*Luminous intensity: I*, candela (cd) The luminous intensity in the perpendicular direction of a surface of  $1/600\ 000\ m^2$  of a black body at the temperature of freezing platinum under a pressure of 101 325 newton per square metre.

Amount of substance: Q, mole (mol) The amount of substance of a system which contains as many elementary entities as there are atoms in 0.012 kg of carbon-12. The elementary entity must be specified and may be an atom, a molecule, an ion, an electron ..., or a specified group of such entities.

# 1.1.2 Supplementary units

*Plane angle:*  $\alpha$ ,  $\beta \varsigma ...$ , radian (rad) The plane angle between two radii of a circle which cut off on the circumference of the circle an arc of length equal to the radius.

Solid angle:  $\Omega$ , steradian (sr) The solid angle which, having its vertex at the centre of a sphere, cuts off an area of the surface of the sphere equal to a square having sides equal to the radius.

# 1.1.3 Notes

*Temperature* At zero K, bodies possess no thermal energy. Specified points (273.16 and 373.16 K) define the Celsius (centigrade) scale (0 and 100°C). In terms of *intervals*,  $1^{\circ}C = 4K$ . In terms of *levels*, a scale Celsius temperature  $\theta_{\infty}$  corresponds to ( $\theta_{\varsigma}$ + 273.16) K.

*Force* The SI unit is the newton (N). A force of 1 N endows a mass of 1 kg with an acceleration of  $1 \text{ m/s}^2$ .

*Weight* The weight of a mass depends on gravitational effect. The standard weight of a mass of 1 kg at the surface of the earth is 9.807 N.

# 1.1.4 Derived units

All physical quantities have units derived from the base and supplementary SI units, and some of them have been given names for convenience in use. A tabulation of those of interest in electrical technology is appended to the list in *Table 1.1*.

Table 1.1	SI base.	supplementary	and and	derived units

Quantity	Unit name	Derivation	Unit symbol
Length	metre		m
Mass	kilogram		kg
Time	second		s
Electric current	ampere		А
Thermodynamic			
temperature	kelvin		K
Luminous			
intensity	candela		cd
Amount of	mole		mol
substance			
Plane angle	radian		rad
Solid angle	steradian	2	sr
Force	newton	kg m/s <sup>2</sup>	Ν
Pressure, stress	pascal	$N/m^2$	Ра
Energy	joule	N m, W s	J
Power	watt	J/s	W
Electric charge,			~
flux	coulomb	A s	С
Magnetic flux	weber	Vs	Wb
Electric potential	volt	J/C	V
Magnetic flux	. 1	xx11 / 2	Ŧ
density	tesla	Wb/m <sup>2</sup>	T
Resistance	ohm	V/A	Ω
Inductance	henry	Wb/A, V s/A	Н
Capacitance	farad	C/V, A s/V	F
Conductance	siemens	A/V s <sup>-1</sup>	S
Frequency	hertz		Hz
Luminous flux	lumen	cd sr	lm 1
Illuminance	lux	$lm/m^2$	lx
Radiation	h a a garanal	$s^{-1}$	Da
activity	becquerel	-	Bq
Absorbed dose	gray	J/kg	Gy
Mass density	kilogram per		$1 ca/m^3$
Dunamia	cubic metre		kg/m³
Dynamic	massal sasand		Pas
viscosity	pascal-second		
Concentration	mole per cubic		mol/ m <sup>3</sup>
Linear velocity	metre per second		
Linear velocity Linear	metre per second		$\frac{m}{s}$
acceleration	metre per second- squared	•	$m/s^2$
Angular velocity	radian per second	I	rad/s
ringular velocity	radian per second	L	cont'd

Table 1.1 (continued)

Quantity	Unit name	Derivation	Unit symbol
Angular acceleration	radian per second- squared		rad/s <sup>2</sup>
Torque	newton metre		Nm
Electric field			
strength	volt per metre		V/m
Magnetic field	-		
strength	ampere per metre		A/m
Current density	ampere per square		,
5	metre		$A/m^2$
Resistivity	ohm metre		Ώm
Conductivity	siemens per metre		S/m
Permeability	henry per metre		H∕m
Permittivity	farad per metre		F/m
Thermal	1		1
capacity	joule per kelvin		J/K
Specific heat	joule per kilogram		- 1
capacity	kelvin		J/(kg K)
Thermal	watt per metre		
conductivity	kelvin		W/(m K)
Luminance	candela per		,()
24111111100	square metre		$cd/m^2$

Decimal multiples and submultiples of SI units are indicated by prefix letters as listed in *Table 1.2*. Thus, kA is the unit symbol for kiloampere, and  $\mu$ F that for microfarad. There is a preference in technology for steps of 10<sup>3</sup>. Prefixes for the kilogram are expressed in terms of the gram: thus, 1000 kg = 1 Mg, not 1 kkg.

Table 1.2 Decimal prefixes

#### 1.1.5 Auxiliary units

Some quantities are still used in special fields (such as vacuum physics, irradiation, etc.) having non-SI units. Some of these are given in *Table 1.3* with their SI equivalents.

# 1.1.6 Conversion factors

Imperial and other non-SI units still in use are listed in *Table 1.4*, expressed in the most convenient multiples or submultiples of the basic SI unit [ ] under classified headings.

# 1.1.7 CGS electrostatic and electromagnetic units

Although obsolescent, electrostatic and electromagnetic units (e.s.u., e.m.u.) appear in older works of reference. Neither system is 'rationalised', nor are the two mutually compatible. In e.s.u. the electric space constant is  $\varepsilon_6 = 1$ , in e.m.u. the magnetic space constant is  $\mu_0 = 1$ ; but the SI units take account of the fact that  $1/\sqrt{(\varepsilon_6\mu_0)}$  is the velocity of electromagnetic wave propagation in free space. *Table 1.5* lists SI units with the equivalent number *n* of e.s.u. and 'ab' ('absolute') for e.m.u. Thus, 1 V corresponds to  $10^{-2}/3$  stV and to  $10^8$  abV, so that 1 stV = 300 V and 1 abV =  $10^{-8}$  V.

# 1.2 Mathematics

Mathematical symbolism is set out in *Table 1.6*. This subsection gives trigonometric and hyperbolic relations, series (including Fourier series for a number of common wave forms), binary enumeration and a list of common derivatives and integrals.

$10^{18}$ exa E $10^{15}$ peta P $10^{12}$ tera T	10 <sup>9</sup> giga G 10 <sup>6</sup> mega M 10 <sup>3</sup> kilo k	$10^{2} hecto h$ $10^{1} deca da$ $10^{-1} deci d$ $10^{-2} centi c$	$10^{-3}$ milli m $10^{-6}$ micro $\mu_{\varsigma}$ $10^{-9}$ nano n	$\frac{10^{-12} \text{ pico p}}{10^{-15} \text{ femto f}}$ $\frac{10^{-18} \text{ atto a}}{10^{-18} \text{ atto a}}$
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Table 1.3 Auxiliary units

Quantity	Symbol	SI		Quantity	Symbol	SI	
Angle				Mass			
degree	(°)	$\pi/180$	rad	tonne	t	1000	kg
minute	Ŏ						U
second	('')			Nucleonics, Radiation			
				becquerel	Bq	1.0	$s^{-1}$
Area				gray	Gy	1.0	J/kg
are	а	100	$m^2$	curie	Ci	$3.7 \times 10^{10}$	Bq
hectare	ha	0.01	km <sup>2</sup>	rad	rd	0.01	Gy
barn	barn	$10^{-28}$	$\frac{km^2}{m^2}$	roentgen	R	$2.6  imes 10^{-4}$	C/kg
Energy				Pressure			, 0
erg	erg	0.1	μJ	bar	b	100	kPa
calorie	cal	4.186	j	torr	Torr	133.3	Ра
electron-volt	eV	0.160	aJ	Time			
gauss-oersted	Ga Oe	7.96	$\mu J/m^3$	minute	min	60	S
Force			• /	hour	h	3600	S
dyne	dyn	10	μN	day	d	86 400	S
Length	2		•	Volume			
Ångstrom	Å	0.1	μm	litre	1 or L	1.0	dm <sup>3</sup>

# Table 1.4 Conversion factors

$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 17.86 \ \text{kg/m} \\ 1.488 \ \text{kg/m} \\ 0.496 \ \text{kg/m} \\ 27.68 \ \text{Mg/m}^3 \\ 16.02 \ \text{kg/m}^3 \\ 1329 \ \text{kg/m}^3 \\ 1329 \ \text{kg/m}^3 \\ 0.1260 \ \text{g/s} \\ 0.2822 \ \text{kg/s} \\ 0.4536 \ \text{kg/s} \\ 7.866 \ \text{cm}^3/\text{s} \\ 1.263 \ \text{cm}^3/\text{s} \\ 1.263 \ \text{cm}^3/\text{s} \\ 1.263 \ \text{cm}^3/\text{s} \\ 4.546 \ \text{dm}^3/\text{s} \\ 10.0 \ \mu\text{N} \\ 9.807 \ \text{N} \\ 0.278 \ \text{N} \\ 4.445 \ \text{N} \\ 9.964 \ \text{kN} \end{array}$
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 0.496 \ \text{kg/m} \\ 27.68 \ \text{Mg/m}^3 \\ 16.02 \ \text{kg/m}^3 \\ 1329 \ \text{kg/m}^3 \\ 1329 \ \text{kg/m}^3 \\ 0.1260 \ \text{g/s} \\ 0.2822 \ \text{kg/s} \\ 0.4536 \ \text{kg/s} \\ 7.866 \ \text{cm}^3/\text{s} \\ 1.263 \ \text{cm}^3/\text{s} \\ 1.263 \ \text{cm}^3/\text{s} \\ 4.546 \ \text{dm}^3/\text{s} \\ \end{array}$
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 27.68 \text{ Mg/m}^3 \\ 16.02 \text{ kg/m}^3 \\ 1329 \text{ kg/m}^3 \\ \end{array}$
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	16.02 kg/m <sup>3</sup> 1329 kg/m <sup>3</sup> 0.1260 g/s 0.2822 kg/s 0.4536 kg/s 7.866 cm <sup>3</sup> /s 0.0283 m <sup>3</sup> /s 1.263 cm <sup>3</sup> /s 75.77 cm <sup>3</sup> /s 4.546 dm <sup>3</sup> /s 10.0 μN 9.807 N 0.278 N 4.445 N
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	1329 kg/m <sup>3</sup> 0.1260 g/s 0.2822 kg/s 0.4536 kg/s 7.866 cm <sup>3</sup> /s 0.0283 m <sup>3</sup> /s 1.263 cm <sup>3</sup> /s 75.77 cm <sup>3</sup> /s 4.546 dm <sup>3</sup> /s 10.0 μN 9.807 N 0.278 N 4.445 N
$\begin{array}{c c c c c c c c c c c c c c c c c c c $	0.1260 g/s 0.2822 kg/s 0.4536 kg/s 7.866 cm <sup>3</sup> /s 0.0283 m <sup>3</sup> /s 1.263 cm <sup>3</sup> /s 75.77 cm <sup>3</sup> /s 4.546 dm <sup>3</sup> /s 10.0 μN 9.807 N 0.278 N 4.445 N
Area $[m^2]$ Flow rate $[kg/s, m^3/s]$ 1 circular mil $506.7  \mu m^2$ 1 tb/h         1 in <sup>2</sup> $645.2  mm^2$ 1 tb/s         1 ft <sup>2</sup> $0.0929  m^2$ 1 ft <sup>3</sup> /h         1 yd <sup>2</sup> $0.8361  m^2$ 1 ft <sup>3</sup> /s         1 acre $4047  m^2$ 1 gal/h         1 mile <sup>2</sup> $2.590  km^2$ 1 gal/min         1 mile <sup>2</sup> $2.590  km^2$ 1 gal/min         1 ft <sup>3</sup> $0.0283  m^3$ 1 dyn         1 yd <sup>3</sup> $0.7646  m^3$ 1 kgf         1 UK gal $4.546  dm^3$ 1 ozf         1 tonf       1 dyn/cm <sup>2</sup> 1 dyn/cm <sup>2</sup> 1 ft/min $5.080  mm/s$ 1 lbf/ft <sup>2</sup> 1 ft/s $0.3048  m/s$ 1 tonf/ft <sup>2</sup> 1 mile/h $0.4470  m/s$ 1 tonf/ft <sup>2</sup> 1 mile/h $0.4470  m/s$ 1 kgf/m <sup>2</sup> 1 knot $0.5144  m/s$ 1 kgf/m <sup>2</sup> 1 keg/s $17.45  mrad/s$ 1 kgf/cm <sup>2</sup>	$\begin{array}{c} 0.2822 \ \text{kg/s} \\ 0.4536 \ \text{kg/s} \\ 7.866 \ \text{cm}^3 \ \text{s} \\ 1.263 \ \text{cm}^3 \ \text{s} \\ 1.263 \ \text{cm}^3 \ \text{s} \\ 75.77 \ \text{cm}^3 \ \text{s} \\ 4.546 \ \text{dm}^3 \ \text{s} \\ \end{array}$
Area $[m^2]$ 1 lb/h         1 circular mil       506.7 $\mu m^2$ 1 ton/h         1 in <sup>2</sup> 645.2 mm <sup>2</sup> 1 lb/s         1 ft <sup>2</sup> 0.0929 m <sup>2</sup> 1 ft <sup>3</sup> /h         1 yd <sup>2</sup> 0.8361 m <sup>2</sup> 1 ft <sup>3</sup> /s         1 acre       4047 m <sup>2</sup> 1 gal/h         1 mile <sup>2</sup> 2.590 km <sup>2</sup> 1 gal/min         1 gal/s       1 gal/s         Volume [m <sup>3</sup> ]       1         1 in <sup>3</sup> 16.39 cm <sup>3</sup> 1 ft <sup>3</sup> 0.0283 m <sup>3</sup> 1 yd <sup>3</sup> 0.7646 m <sup>3</sup> 1 UK gal       4.546 dm <sup>3</sup> 1 bf       1 conf         Acceleration [m/s <sup>2</sup> , rad/s <sup>2</sup> ]       1 dyn/cm <sup>2</sup> 1 ft/s       0.3048 m/s       1 lbf/in <sup>2</sup> 1 in/s       25.40 mm/s       1 lbf/in <sup>2</sup> 1 ft/s       0.3048 m/s       1 conf/ft <sup>2</sup> 1 knot       0.5144 m/s       1 kgf/m <sup>2</sup> 1 knot       0.5144 m/s       1 kgf/cm <sup>2</sup>	$\begin{array}{c} 0.2822 \ \text{kg/s} \\ 0.4536 \ \text{kg/s} \\ 7.866 \ \text{cm}^3/\text{s} \\ 0.0283 \ \text{m}^3/\text{s} \\ 1.263 \ \text{cm}^3/\text{s} \\ 75.77 \ \text{cm}^3/\text{s} \\ 4.546 \ \text{dm}^3/\text{s} \\ \end{array}$
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	0.4536 kg/s 7.866 cm <sup>3</sup> /s 0.0283 m <sup>3</sup> /s 1.263 cm <sup>3</sup> /s 75.77 cm <sup>3</sup> /s 4.546 dm <sup>3</sup> /s 10.0 μN 9.807 N 0.278 N 4.445 N
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	7.866 cm <sup>3</sup> /s 0.0283 m <sup>3</sup> /s 1.263 cm <sup>3</sup> /s 75.77 cm <sup>3</sup> /s 4.546 dm <sup>3</sup> /s 10.0 μN 9.807 N 0.278 N 4.445 N
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	0.0283 m <sup>3</sup> /s 1.263 cm <sup>3</sup> /s 75.77 cm <sup>3</sup> /s 4.546 dm <sup>3</sup> /s 10.0 μN 9.807 N 0.278 N 4.445 N
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	1.263 cm <sup>3</sup> /s 75.77 cm <sup>3</sup> /s 4.546 dm <sup>3</sup> /s 10.0 μN 9.807 N 0.278 N 4.445 N
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	75.77 cm <sup>3</sup> /s 4.546 dm <sup>3</sup> /s 10.0 μN 9.807 N 0.278 N 4.445 N
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	4.546 dm <sup>3</sup> /s 10.0 μN 9.807 N 0.278 N 4.445 N
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	9.807 N 0.278 N 4.445 N
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$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	0.278 N 4.445 N
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	4.445 N
$\begin{array}{cccc} Velocity \ [m/s, rad/s] & 1 \ tonf \\ Acceleration \ [m/s^2, rad/s^2] & 1 \ dyn/cm^2 \\ 1 \ ft/min & 5.080 \ mm/s & 1 \ lbf/ft^2 \\ 1 \ in/s & 25.40 \ mm/s & 1 \ lbf/in^2 \\ 1 \ ft/s & 0.3048 \ m/s & 1 \ tonf/tr^2 \\ 1 \ mile/h & 0.4470 \ m/s & 1 \ tonf/in^2 \\ 1 \ knot & 0.5144 \ m/s & 1 \ kgf/m^2 \\ 1 \ deg/s & 17.45 \ mrad/s & 1 \ kgf/cm^2 \end{array}$	
Acceleration $[m/s^2, rad/s^2]$ 1 dyn/cm²         1 ft/min       5.080 mm/s       1 lbf/ft²         1 in/s       25.40 mm/s       1 lbf/in²         1 ft/s       0.3048 m/s       1 tonf/ft²         1 mile/h       0.4470 m/s       1 tonf/in²         1 knot       0.5144 m/s       1 kgf/m²         1 deg/s       17.45 mrad/s       1 kgf/cm²	
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	0.10 Pa
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	47.88 Pa
$\begin{array}{ccccccc} 1 & \text{mile/h} & 0.4470  \text{m/s} & 1  \text{tonf/in}^2 \\ 1 & \text{knot} & 0.5144  \text{m/s} & 1  \text{kgf/m}^2 \\ 1 & \text{deg/s} & 17.45  \text{mrad/s} & 1  \text{kgf/cm}^2 \end{array}$	6.895 kPa
$\begin{array}{ccccccc} 1 \ knot & 0.5144 \ m/s & 1 \ kgf/m^2 \\ 1 \ deg/s & 17.45 \ mrad/s & 1 \ kgf/cm^2 \end{array}$	107.2 kPa
$1 \text{ deg/s}$ $17.45 \text{ mrad/s}$ $1 \text{ kgf/cm}^2$	15.44 MPa 9.807 Pa
	98.07 kPa
	133.3 Pa
1  rev/s $6.283  rad/s$ $1  inHg$	3.386 kPa
$1 \text{ ft/s}^2$ 0.3048 m/s <sup>2</sup> 1 inH <sub>2</sub> O	149.1 Pa
1 mile/h per s $0.4470 \mathrm{m/s^2}$ 1 ftH <sub>2</sub> O	2.989 kPa
Mass [kg] Torque [Nm]	
1 oz 28.35 g 1 ozf in	7.062 nN m
1 lb 0.454 kg 1 lbf in	0.113 N m
1 slug 14.59 kg 1 lbf ft	1.356 N m
1 cwt 50.80 kg 1 tonf ft 1 UKton 1016 kg 1 kgf m	3.307 kN m 9.806 N m
	9.000 IN III
Energy [J], Power [W]Inertia $[kg m^2]$ 1 ft lbf1.356 JMomentum $[kg m/s, kg m^2/s]$	
$1 \text{ m kgf} \qquad 9.807 \text{ J} \qquad 1 \text{ oz in}^2$	$0.018 \mathrm{g}\mathrm{m}^2$
$1 \text{ Btu} \qquad 1055 \text{ J} \qquad 1 \text{ lb in}^2$	$0.293 \mathrm{g}\mathrm{m}^2$
1 therm $105.5 \text{ kJ}$ $1 \text{ lb ft}^2$	$0.0421 \mathrm{kg}\mathrm{m}^2$
1 hp h $2.685 \text{ MJ}$ 1 slug ft <sup>2</sup>	$1.355 \mathrm{kg  m^2}$
$1 \text{ kW h} \qquad 3.60 \text{ MJ} \qquad 1 \text{ ton } \text{ft}^2$	$94.30 \mathrm{kg}\mathrm{m}^2$
1 Btu/h $0.293 \text{ W}$ 1 lb ft/s	0.138  kg m/s
$1 \text{ ft lbf/s}$ $1.356 \text{ W}$ $1 \text{ lb ft}^2/\text{s}$	$0.042  \text{kg}  \text{m}^2/\text{s}$
1 m kgf/s 9.807 W 1 hp 745.9 W Viscosity [Pa s, m <sup>2</sup> /s]	
l np l poise	9.807 Pa s
Thermal quantities [W, J, kg, K] 1 kgf s/m <sup>2</sup>	9.807 Pa s
$1 \text{ W/in}^2$ $1.550 \text{ kW/m}^2$ $1 \text{ lbf s/ft}^2$	47.88 Pa s
$1 \text{ Btu}/(\text{ft}^2 \text{ h})$ $3.155 \text{ W/m}^2$ $1 \text{ lbf h/ft}^2$	172.4 kPa s
$1 \text{ Btu/(ft^3 h)}$ $10.35 \text{ W/m}^3$ $1 \text{ stokes}$	$1.0  \text{cm}^2/\text{s}$
$1 \text{ Btu}/(\text{ft h}^{\circ}\text{F})$ $1.731 \text{ W}/(\text{m K})$ $1 \frac{\text{in}^2/\text{s}}{1.62^2/\text{s}}$	$6.452 \text{ cm}^2/\text{s}$
$1 \text{ ft } \text{lbf/lb}$ 2.989 J/kg $1 \text{ ft}^2/\text{s}$	$929.0  \text{cm}^2/\text{s}$
1 Btu/lb         2326 J/kg           1 Btu/ft <sup>3</sup> 37.26 KJ/m <sup>3</sup> Illumination [cd, lm]	
1 ft lbf/(lb °F) $5.380 \text{ J/(kg K)}$ $1 \text{ lm/ft}^2$	$10.76  \text{lm/m}^2$
$1 \text{ Btu/(lb °F)}$ $4.187 \text{ kJ/(kg K)}$ $1 \text{ cd/ft}^2$	$10.76 \text{ cd/m}^2$
$1 \text{ Btu/(ft}^3 ^{\circ}\text{F})$ 67.07 kJ/m <sup>3</sup> K $1 \text{ cd/in}^2$	$1550  cd/m^2$

Table 1.5	Relation	between	SI,	e.s.	and	e.m.	units
-----------	----------	---------	-----	------	-----	------	-------

Quantity	SI unit		Equivalent	number n of	
		<i>e.s</i>	.u.	е.т.	<i>u</i> .
Length	m	10 <sup>2</sup>	cm	10 <sup>2</sup>	cm
Mass	kg	$10^{3}$	g	$10^{3}$	g
Time	s	1	s	1	s
Force	N	$10^{5}$	dyn	$10^{5}$	dyn
Torque	Nm	$10^{7}$	dyn cm	107	dyn cm
Energy	J	$10^{7}$	erg	107	erg
Power	Ŵ	$10^{7}$	erg/s	107	erg/s
Charge, electric flux	Ĉ	$3 \times 10^{9}$	stC	$10^{-1}$	abC
density	$\tilde{C}/m^2$	$3 \times 10^5$	stC/cm <sup>2</sup>	$10^{-5}$	abC/cm <sup>2</sup>
Potential, e.m.f.	V	$10^{-2}/3$	stV	$10^{8}$	abV
Electric field strength	V/m	$10^{-4}/3$	stV/cm	$10^{6}$	abV/cm
Current	A	$3 \times 10^{9}$	stA	$10^{-1}$	abA
density	$A/m^2$	$3 \times 10^{5}$	stA/cm <sup>2</sup>	$10^{-5}$	abA/cm <sup>2</sup>
Magnetic flux	Wb	$10^{-2}/3$	stWb	$10^{8}$	Mx
density	Т	$10^{-6}/3$	stWb/cm <sup>2</sup>	$10^{4}$	Gs
Mag. fd. strength	Ā/m	$12\pi \times 10^7$	stA/cm	$4\pi \times 10^{-3}$	Ōe
M.M.F.	A	$12\pi \times 10^9$	stA	$4\pi  imes 10^{-1}$	Gb
Resistivity	Ωm	$10^{-9}/9$	st $\Omega$ cm	10 <sup>11</sup>	abΩ cm
Conductivity	S/m	$9 \times 10^{9}$	stS/cm	$10^{-11}$	abS/cm
Permeability (abs)	H/m	$10^{-13}/36\pi\varsigma$		$10^{7}/4\pi\varsigma$	
Permittivity (abs)	F/m	$36\pi \times 10^9$		$4\pi  imes 10^{-11}$	
Resistance	$\Omega^{'}$	$10^{-11}/9$	$st\Omega$	109	abΩ
Conductance	S	$9 \times 10^{11}$	stS	$10^{-9}$	abS
Inductance	Н	$10^{-12}/9$	stH	$10^9$	cm
Capacitance	F	$9 \times 10^{11}$	cm	$9 \times 10^{11}$	abF
Reluctance	A/Wb	$36\pi \times 10^{11}$		$4\pi  imes 10^{-8}$	Gb/Mx
Permeance	Wb/A	$10^{11}/36\pi\varsigma$		$10^9/4\pi\varsigma$	Mx/Gb

Gb = gilbert; Gs = gauss; Mx = maxwell; Oe = oersted.

# 1.2.1 Trigonometric relations

The trigonometric functions (sine, cosine, tangent, cosecant, secant, cotangent) of an angle  $\theta$  are based on the circle, given by  $x^2 + y^2 = h^2$ . Let two radii of the circle enclose an angle  $\theta_{\varsigma}$  and form the sector area  $S_c = (\pi h^2)(\theta/2\pi)$  shown shaded in *Figure 1.1* (left): then  $\theta_{\varsigma}$ can be defined as  $2S_c/h^2$ . The *right-angled* triangle with sides *h* (hypotenuse), *a* (adjacent side) and *p* (opposite side) give ratios defining the trigonometric functions

$\sin\theta = p/h$	$\operatorname{cosec}\theta = 1/\sin\theta = h/p$
$\cos\theta = a/h$	sec $\theta = 1/\cos\theta = h/a$
$\tan \theta = p/a$	$\cot a \theta = 1/\tan \theta = a/p$

In *any* triangle (*Figure 1.1*, right) with angles, *A*, *B* and *C* at the corners opposite, respectively, to sides *a*, *b* and *c*, then  $A+B+C=\pi$  rad (180°) and the following relations hold:

 $a = b \cos C + c \cos B$   $b = c \cos A + a \cos C$   $c = a \cos B + b \cos A$   $a/\sin A = b/\sin B = c/\sin C$   $a^2 = b^2 + c^2 + 2bc \cos A$  $(a+b)/(a-b) = (\sin A + \sin B)/(\sin A - \sin B) \Leftarrow$ 

Other useful relationships are:

 $sin(x \pm y) = sin x \cdot cos y \pm cos x \cdot sin y$  $cos(x \pm y) = cos x \cdot cos y \mp sin x \cdot sin y$ 

 $\tan(x \pm y) = (\tan x \cdot \tan y)/(1 \mp \tan x \cdot \tan y) \Leftrightarrow$   $\sin^2 x = \frac{1}{2}(1 - \cos 2x)\cos^2 x = -\frac{1}{2}(1 + \cos 2x)$   $\sin^2 x + \cos^2 x = 1 \quad \sin^3 x = -\frac{1}{4}(3\sin x - \sin 3x)$  $\cos^3 x = \frac{1}{4}(3\cos x + \cos 3x)$ 

$$\sin x \pm \sin y = 2 \begin{bmatrix} \cos_{\frac{1}{2}}(x-y) \cdot \frac{\sin_{\frac{1}{2}}(x+y)}{\cos^{\frac{1}{2}}(x+y)} \end{bmatrix}$$
$$\cos x \pm \cos y = -2 \begin{bmatrix} \cos_{\frac{1}{2}}(x-y) \cdot \frac{\sin_{\frac{1}{2}}(x+y)}{\cos^{\frac{1}{2}}(x+y)} \end{bmatrix}$$

$$\tan x \pm \tan y = \sin(x \pm y)/\cos x \cdot \cos y$$
  

$$\sin^2 x - \sin^2 y = \sin(x + y) \cdot \sin(x - y) \Leftarrow$$
  

$$\cos^2 x - \cos^2 y = -\sin(x + y) \cdot \sin(x - y) \Leftarrow$$
  

$$\cos^2 x - \sin^2 y = \cos(x + y) \cdot \cos(x - y) \Leftarrow$$

$$d(\sin x)/dx = \cos x \qquad \qquad \int \sin x \cdot dx = -\cos x + k$$
  
$$d(\cos x)/dx = -\sin x \qquad \int \cos x \cdot dx = \sin x + k$$
  
$$d(\tan x)/dx = \sec^2 x \qquad \int \tan x \cdot dx = -\ln|\cos x| + k$$

Values of sin  $\theta$ , cos  $\theta$  and tan  $\theta$  for  $0^{\circ \leq} \theta < 90^{\circ}$  (for  $0 < \theta \varsigma < 1.571$  rad) are given in *Table 1.7* as a check list, as they can generally be obtained directly from calculators.

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#### Table 1.6 Mathematical symbolism

Term	Symbol
Base of natural logarithms	e (=2.718 28)
Complex number	$\mathbf{C} = A + \mathbf{j}B = C \exp(\mathbf{j}\theta)$
argument; modulus	$= C \ \angle \ \theta_{\varsigma}$ arg <b>C</b> = $\theta$ ; mod <b>C</b> = C
conjugate	$\mathbf{C}^{\mathbf{*}} = A - \mathbf{j}B = C \exp(-\mathbf{j}\theta)$ $= C \ \angle -\theta\varsigma$
real part; imaginary part Co-ordinates	Re $C = A$ ; Im $C = B$
cartesian	<i>x</i> , <i>y</i> , <i>z</i>
cylindrical; spherical	$r, \phi, z; r, \theta, \phi\varsigma$
Function of x	
general	f(x), g(x), F(x)
Bessel	$J_n(x)$
circular	$\sin x, \cos x, \tan x \dots$
inverse	$\arcsin x, \arccos x,$
1:66	$\arctan x \dots$
differential	d <i>x</i>
partial	$\partial x$
exponential	$\exp(x)$
hyperbolic	$\sinh x$ , $\cosh x$ , $\tanh x \dots$
inverse	arsinh x, arcosh x, artanh $x \dots$
increment	$\Delta x,  \delta x$
limit	lim x
logarithm	
base b	log <sub>b</sub> x
common; natural	$lg x; ln x (or log x; log_e x)$
Matrix	A, B
complex conjugate	A*, B*
product	AB
square, determinant	det A
inverse	$A^{-1}$
transpose	$A^{t}$
unit	Ι
Operator	·· (- (1/1))
Heaviside	$p (\equiv \langle \mathbf{d} / \mathbf{d} t \rangle)$
impulse function	$\delta(t)$
Laplace $L[f(t)] = F(s)$	$s (= \sigma + j\omega)$ $\nabla \Leftarrow$
nabla, del rotation $\pi/2$ rad:	
rotation $\pi/2$ rad; $2\pi/3$ rad	J h
$2\pi/3$ rad	
step function Vector	H(t), u(t)
curl of A	A, a, B, b curl A $\nabla \times A$
divergence of A	curl $\mathbf{A}, \nabla \times \mathbf{A}$ div $\mathbf{A}, \nabla \cdot \mathbf{A}$
gradient of $\phi_{\varsigma}$	grad $\phi$ , $\nabla \not\Leftrightarrow \varsigma$
product: scalar; vector	$\mathbf{A} \cdot \mathbf{B}; \mathbf{A} \times \mathbf{B}$

Table 1.7	Trigonometric	functions of $\theta_{\varsigma}$
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$\theta\varsigma$		sin $\theta\varsigma$	$\cos \theta \varsigma$	tan $\theta\varsigma$	
deg	rad				
0	0.0	0.0	1.0	0.0	
5	0.087	0.087	0.996	0.087	
10	0.175	0.174	0.985	0.176	
15	0.262	0.259	0.966	0.268	
20	0.349	0.342	0.940	0.364	
25	0.436	0.423	0.906	0.466	
30	0.524	0.500	0.866	0.577	
35	0.611	0.574	0.819	0.700	
40	0.698	0.643	0.766	0.839	
45	0.766	0.707	0.707	1.0	
50	0.873	0.766	0.643	1.192	
55	0.960	0.819	0.574	1.428	
60	1.047	0.866	0.500	1.732	
65	1.134	0.906	0.423	2.145	
70	1.222	0.940	0.342	2.747	
75	1.309	0.966	0.259	3.732	
80	1.396	0.985	0.174	5.671	
85	1.484	0.996	0.097	11.43	
90	1.571	1.0	0.0	$\infty \Leftarrow$	

# 1.2.2 Exponential and hyperbolic relations

Exponential functions For a positive datum ('real') number u, the exponential functions exp(u) and exp(-u)are given by the summation to infinity of the series

$$\exp(\pm u) = \mathbf{k} \pm u + u^2/2! \pm u^3/3! + u^4/4! \pm \cdots \neq$$

with  $\exp(+u)$  increasing and  $\exp(-u)$  decreasing at a rate proportional to *u*.

If u = 1, then

$$exp(+1) = 1 + 1 + 1/2 + 1/6 + 1/24 + \dots = e = 2.718 \dots \Leftrightarrow exp(-1) = 1 - 1 + 1/2 - 1/6 + 1/24 - \dots = 1/e = 0.368 \dots \Leftrightarrow$$

In the electrical technology of transients, u is most commonly a negative function of time t given by u = -(t/T). It then has the graphical form shown in Figure 1.2 (left) as a time dependent variable. With an initial value k, i.e.  $y = k \exp(-t/T)$ , the rate of reduction with time is  $dy/dt = \Leftarrow$  $-(k/T)\exp(-t/T)$ . The initial rate at t=0 is -k/T. If this rate were maintained, y would reach zero at t = T, defining the time constant T. Actually, after time T the value of y is  $\bar{k}$  $\exp(-t/T) = k \exp(-1) = 0.368k$ . Each successive interval T decreases y by the factor 0.368. At a time t = 4.6T the value of y is 0.01k, and at t = 6.9T it is 0.001k.

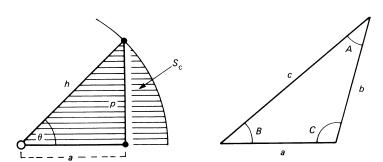
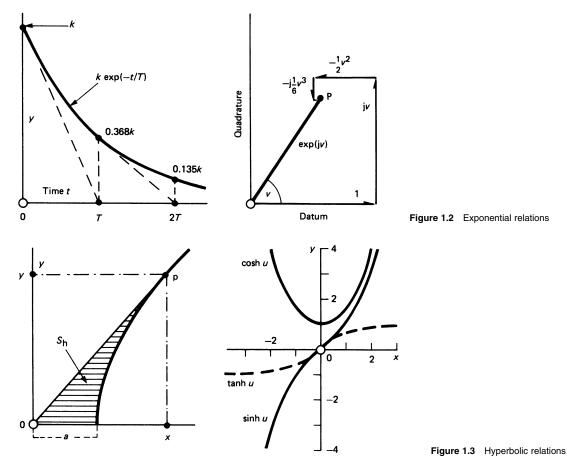


Figure 1.1 Trigonometric relations



If *u* is a quadrature ('imaginary') number  $\pm jv$ , then

$$\exp(\pm jv) = 1 \pm jv - v^2/2! \mp jv^3/3! + v^4/4! \pm \dots$$

because  $j^2 = -1$ ,  $j^3 = -j1$ ,  $j^4 = +1$ , etc. Figure 1.2 (right) shows the summation of the first five terms for exp(j1), i.e.

$$\exp(j1) = 1 + j1 - 1/2 - j1/6 + 1/24$$

a complex or expression converging to a point P. The length OP is unity and the angle of OP to the datum axis is, in fact, 1 rad. In general,  $\exp(jv)$  is equivalent to a shift by  $\angle v$  rad. It follows that  $\exp(\pm jv) = \cos v \pm j \sin v$ , and that

$$\exp(jv) + \exp(-jv) = 2\cos v \quad \exp(jv) - \exp(-jv) = j2\sin v$$

For a complex number (u + jv), then

$$\exp(u + \mathbf{j}v) = \exp(u) \cdot \exp(\mathbf{j}v) = \exp(u) \cdot \angle v$$

Hyperbolic functions A point P on a rectangular hyperbola  $(x_i a)^2 \rightarrow \notin y_i a)^2 = 1$  defines the hyperbolic 'sector' area  $S_h = \frac{1}{2}a^2 \ln[(x/a - (y/a)]$  shown shaded in Figure 1.3 (left). By analogy with  $\theta \in 2S_c/h^2$  for the trigonometrical angle  $\theta$ , the hyperbolic entity (not an angle in the ordinary sense) is  $u = 2S_h/a^2$ , where a is the major semi-axis. Then the hyperbolic functions of u for point P are:

$$sinh u = y/a$$
 $cosech u = a/y$  $cosh u = x/a$  $sech u = a/x$  $tanh u = y/x$  $coth u = x/y$ 

The principal relations yield the curves shown in the diagram (right) for values of *u* between 0 and 3. For higher values  $\sinh u$  approaches  $\pm \cosh u$ , and  $\tanh u$  becomes asymptotic to  $\pm 1$ . Inspection shows that  $\cosh(-u) = \cosh u$ ,  $\sinh(-u) = -\sinh u$  and  $\cosh^2 u - \sinh^2 u = 1$ .

The hyperbolic functions can also be expressed in the exponential form through the series

$$\cosh u = 1 + u^2/2! + u^4/4! + u^6/6! + \dots \Leftarrow$$
  
$$\sinh u = u + u^3/3! + u^5/5! + u^7/7! + \dots \Leftarrow$$

so that

$$\cosh u = \frac{1}{2} [\exp(u) + \exp(-u)] \Leftrightarrow \sinh u = \frac{1}{2} [\exp(u) - \exp(-u)]$$
  
$$\cosh u + \sinh u = \exp(u) \Leftrightarrow \cosh u - \sinh u = \exp(-u) \Leftrightarrow$$

Other relations are:

$$\sinh u + \sinh v = 2 \sinh \frac{1}{2}(u+v) \cdot \cosh \frac{1}{2}(u-v)$$
  

$$\cosh u + \cosh v = 2 \cosh \frac{1}{2}(u+v) \cdot \cosh \frac{1}{2}(u-v)$$
  

$$\cosh u - \cosh v = 2 \sinh \frac{1}{2}(u+v) \cdot \sinh \frac{1}{2}(u-v)$$
  

$$\sinh(u\pm v) = \sinh u \cdot \cosh v \pm \cosh u \cdot \sinh v$$
  

$$\cosh(u\pm v) = \cosh u \cdot \cosh v \pm \sinh u \cdot \sinh v$$
  

$$\tanh(u\pm v) = (\tanh u \pm \tanh v)/(1 \pm \tanh u \cdot \tanh v) \Leftarrow$$

Table 1.8 Exponential and hyperbolic functions

и	exp(u)	exp(-u)	sinh u	cosh u	tanh u
0.0	1.0	1.0	0.0	1.0	0.0
0.1	1.1052	0.9048	0.1092	1.0050	0.0997
0.2	1.2214	0.8187	0.2013	1.0201	0.1974
0.3	1.3499	0.7408	0.3045	1.0453	0.2913
0.4	1.4918	0.6703	0.4108	1.0811	0.3799
0.5	1.6487	0.6065	0.5211	1.1276	0.4621
0.6	1.8221	0.5488	0.6367	1.1855	0.5370
0.7	2.0138	0.4966	0.7586	1.2552	0.6044
0.8	2.2255	0.4493	0.8881	1.3374	0.6640
0.9	2.4596	0.4066	1.0265	1.4331	0.7163
1.0	2.7183	0.3679	1.1752	1.5431	0.7616
1.2	3.320	0.3012	1.5095	1.8107	0.8337
1.4	4.055	0.2466	1.9043	2.1509	0.8854
1.6	4.953	0.2019	2.376	2.577	0.9217
1.8	6.050	0.1653	2.942	3.107	0.9468
2.0	7.389	0.1353	3.627	3.762	0.9640
2.303	10.00	0.100	4.950	5.049	0.9802
2.5	12.18	0.0821	6.050	6.132	0.9866
2.75	15.64	0.0639	7.789	7.853	0.9919
3.0	20.09	0.0498	10.02	10.07	0.9951
3.5	33.12	0.0302	16.54	16.57	0.9982
4.0	54.60	0.0183	27.29	27.31	0.9993
4.5	90.02	0.0111	45.00	45.01	0.9998
4.605	100.0	0.0100	49.77	49.80	0.9999
5.0	148.4	0.0067	74.20	74.21	0.9999
5.5	244.7	0.0041	122.3	(¢osh u	(tanh u
6.0	403.4	0.0025	201.7	$l = \sinh u$	l = 1.0
6.908	1000	0.0010	500	$=\frac{1}{2}\exp(u)$	

 $\begin{aligned} \sinh(u \pm jv) &= (\sinh u \cdot \cos v) \pm j(\cosh u \cdot \sin v) \Leftarrow \\ \cosh(u \pm jv) &= (\cosh u \cdot \cos v) \pm j(\sinh u \cdot \sin v) \Leftarrow \\ d(\sinh u)/du &= \cosh u \qquad \int \sinh u \cdot du = \cosh u \\ d(\cosh u)/du &= \sinh u \qquad \int \cosh u \cdot du = \sinh u \end{aligned}$ 

Exponential and hyperbolic functions of *u* between zero and 6.908 are listed in *Table 1.8*. Many calculators can give such values directly.

# 1.2.3 Bessel functions

Problems in a wide range of technology (e.g. in eddy currents, frequency modulation, etc.) can be set in the form of the Bessel equation

 $\frac{\mathrm{d}^2 y}{\mathrm{d} x^2} + \frac{1}{x} \cdot \frac{\mathrm{d} y}{\mathrm{d} x} + \left[1 - \frac{n^2}{x^2}\right] y = \mathbf{0}$ 

and its solutions are called Bessel functions of order n. For n = 0 the solution is

$$J_0(x) = 4 - (x^2/2^2) + (x^4/2^2 \cdot 4^2) - (x^6/2^2 \cdot 4^2 \cdot 6^2) + \dots \Leftarrow$$

and for n = 1, 2, 3...

$$J_n(x) = \underbrace{\frac{x^n}{2^n n!}}_{n!} \left[ 1 - \underbrace{\frac{x^2}{2(2n+2)}}_{\Leftarrow} + \underbrace{\frac{x^4}{2 \cdot 4(2n+2)(2n+4)}}_{\leftarrow} \cdots \right]_{=}^{+}$$

*Table 1.9* gives values of  $J_n(x)$  for various values of *n* and *x*.

# 1.2.4 Series

*Factorials* In several of the following the factorial (n!) of integral numbers appears. For *n* between 2 and 10 these are

$2! = \Leftarrow 2$ $3! = \Leftarrow 6$ $4! = \Leftarrow 24$ $5! = \Leftarrow 120$ $6! = \Leftarrow 720$ $7! = \Leftarrow 5040$ $8! = \Leftarrow 40320$ $9! = \Leftarrow 362880$	$1/2! = 0.5$ $1/3! = 0.1667$ $1/4! = 0.417 \times 10^{-1}$ $1/5! = 0.833 \times 10^{-2}$ $1/6! = 0.139 \times 10^{-2}$ $1/7! = 0.198 \times 10^{-3}$ $1/8! = 0.248 \times 10^{-4}$ $1/9! = 0.276 \times 10^{-5}$
9! = 40320 9! = 40320 10! = 3628800	$1/0! = 0.248 \times 10^{-5}$ $1/9! = 0.276 \times 10^{-5}$ $1/10! = 0.276 \times 10^{-6}$

Progression

Arithmetic 
$$a + (a + d) + (a + 2d) + \dots + [a + (n - 1)d]$$
  
=  $\frac{1}{2}n$  (sum of 1st and *n*th terms)

Geometric 
$$a + ar + ar^2 + \dots + ar^{n-1} = a(1-r^n)/(1-r)$$

Trigonometric See Section 1.2.1. Exponential and hyperbolic See Section 1.2.2. *Binomial* 

$$(1 \pm x)^{n} = 4 \pm nx + 4 = \frac{n(n-1)}{2!} \div x^{2} \pm \frac{n(n-1)(n-2)}{3!} \div x^{3} + \dots \Leftrightarrow$$
$$+ (-1)^{r} \frac{n!}{r!(n-r)!} x^{r} + \dots \Leftrightarrow$$
$$(a \pm x)^{n} = 4^{n} [1 \pm (x/a)]^{n}$$

Binomial coeffi	cients n!/[r!	(n-r)!]	are tabulated:
-----------------	---------------	---------	----------------

Term $r =$	=⇔0	1	2	3	4	5	6	7	8	9	10
n = 1	1	1									
2	1	2	1								
3	1	3	3	1							
4	1	4	6	4	1						
5	1	5	10	10	5	1					
6	1	6	15	20	15	6	1				
7	1	7	21	35	35	21	7	1			
8	1	8	28	56	70	56	28	8	1		
9	1	9	36	84	126	126	84	36	9	1	
10	1	10	45	120	210	252	210	120	45	10	1

*Power* If there is a power series for a function f(h), it is given by

$$f(h) = f(0) + hf^{(i)}(0) + (h^2/2!)f^{(ii)}(0) + (h^3/3!)f^{(iii)}(0) + \dots \Leftrightarrow$$
$$+ (h^r/r!)f^{(r)}(0) + \dots \Leftrightarrow \qquad (Maclaurin) \Leftrightarrow$$

$$f(x+h) = f(x) + hf^{(i)}(x) + (h^2/2!)f^{(ii)}(x) + \dots \Leftrightarrow$$
$$+ (h^r/r!)f^{(r)}(x) + \dots \Leftrightarrow (Taylor) \Leftrightarrow$$

Permutation, combination

$${}^{n}P_{r} = n(n-1)(n-2)(n-3)\dots(n-r+1) = n!/(n-r)!$$
  
$${}^{n}C_{r} = (1/r!)[n(n-1)(n-2)(n-3)\dots(n-r+1)] = n!/r!(n-r)!$$
  
Bessel See Section 1.2.3.  
Fourier See Section 1.2.5.

# 1.2.5 Fourier series

A univalued periodic wave form  $f(\theta)$  of period  $2\pi \varsigma$  is represented by a summation in general of sine and cosine waves of fundamental period  $2\pi \varsigma$  and of integral harmonic orders n (=2, 3, 4, ...) as

$$f(\theta) = c_0 + a_1 \cos \theta + a_2 \cos 2\theta + \dots + a_n \cos n\theta + \dots \neq b_1 \sin \theta + b_2 \sin 2\theta + \dots + b_n \sin n\theta + \dots \neq b_n \sin n\theta + \dots \neq b_n \sin n\theta + \dots \neq b_n \sin n\theta + \dots = b$$

The mean value of  $f(\theta)$  over a full period  $2\pi qs$ 

$$c_0 = \frac{1}{2\pi\varsigma} \int_0^{2\pi\varsigma} f(\theta) \cdot \mathrm{d}\theta\varsigma$$

and the harmonic-component amplitudes a and b are

$$a_n = \underset{\pi \varsigma}{\overset{1}{\Leftarrow}} \int_{0}^{2\pi\varsigma} f(\theta) \cdot \cos n\theta\varsigma \, \mathrm{d}\theta, \varsigma \quad b_n = \underset{\pi \varsigma}{\overset{1}{\Leftarrow}} \int_{0}^{2\pi\varsigma} f(\theta) \cdot \sin n\theta\varsigma \, \mathrm{d}\theta\varsigma$$

Table 1.10 gives for a number of typical wave forms the harmonic series in square brackets, preceded by the mean value  $c_0$  where it is not zero.

# 1.2.6 Derivatives and integrals

Some basic forms are listed in *Table 1.11*. Entries in a given column are the integrals of those in the column to its left and the derivatives of those to its right. Constants of integration are omitted.

# 1.2.7 Laplace transforms

Laplace transformation is a method of deriving the response of a system to any stimulus. The system has a basic equation of behaviour, and the stimulus is a pulse, step, sine wave or other variable with time. Such a response involves integration: the Laplace transform method removes integration difficulties, as tables are available for the direct solution of a great variety of problems. The process is analogous to evaluation (for example) of  $y = 2.1^{3.6}$  by *transformation* into a logarithmic form log  $y = 3.6 \times \log(2.1)$ , and a subsequent *inverse transformation* back into arithmetic by use of a table of antilogarithms.

The Laplace transform (L.t.) of a time-varying function f(t) is

$$\mathbf{L}[f(t)] = \mathbf{F}(\mathbf{s}) = \oint_{\mathbf{q}}^{\infty} \exp(-st) \cdot f(t) \cdot \mathbf{d}t$$

and the inverse transformation of F(s) to give f(t) is

$$\mathbf{L}^{-1}[\mathbf{F}(\mathbf{s})] = f(t) = \lim \frac{1}{2\pi} \int_{q-j\omega\varsigma}^{q+j\omega\varsigma} \exp(st) \cdot \mathbf{F}(\mathbf{s}) \cdot \mathbf{ds}$$

The process, illustrated by the response of a current i(t) in an electrical network of impedance z to a voltage v(t)applied at t = 0, is to write down the transform equation

$$I(s) = V(s)/Z(s) \Leftarrow$$

where I(s) is the L.t. of the current i(t), V(s) is the L.t. of the voltage v(t), and Z(s) is the *operational* impedance. Z(s) is obtained from the network resistance R, inductance L and capacitance C by leaving R unchanged but replacing L by Ls and C by 1/Cs. The process is equivalent to writing the network impedance for a steady state frequency  $\omega_{cand}$  then replacing j $\omega_{c}$ by s. V(s) and Z(s) are polynomials in s: the quotient V(s)/Z(s) is reduced algebraically to a form recognisable in the transform table. The resulting current/time relation i(t) is read out: it contains the complete solution. However, if at t=0 the network has initial energy (i.e. if currents flow in inductors or charges are stored in capacitors), the equation becomes

$$I(s) = [V(s) + U(s)]/Z(s) \Leftarrow$$

where U(s) contains such terms as  $LI_0$  and  $(1/s)V_0$  for the inductors or capacitors at t = 0.

A number of useful transform pairs is listed in Table 1.12.

#### **1.2.8 Binary numeration**

A number N in decimal notation can be represented by an ordered set of binary digits  $a_n, a_{n-2}, \ldots, a_2, a_1, a_0$  such that

$$N = 2^{n}a_{n} + 2^{n-1}a_{n-1} + \dots + 2a_{1} + a_{0}$$

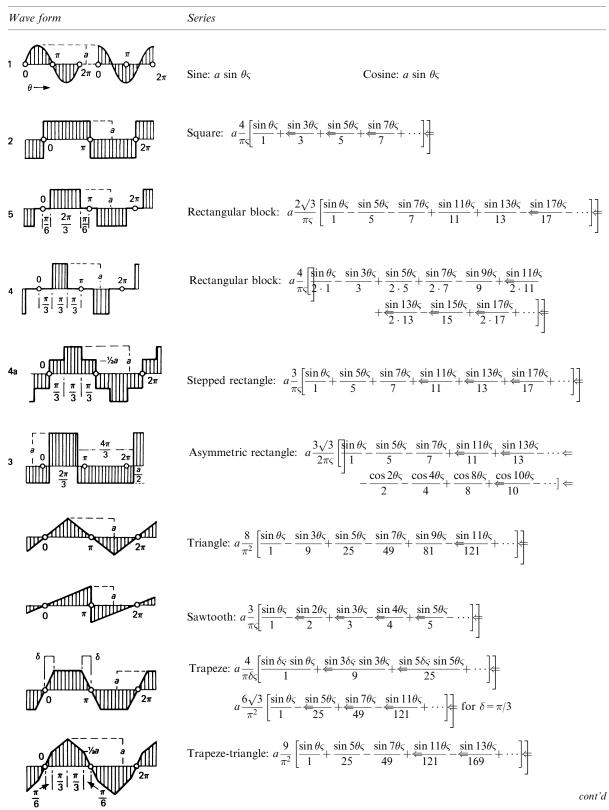
Decimal	1	2	3	4	5	6	7	8	9	10	100
Binary	1	10	11	100	101	110	111	1000	1001	1010	1100100

Tab	le 1	1.9	Bessel	functions	$J_n(x)$
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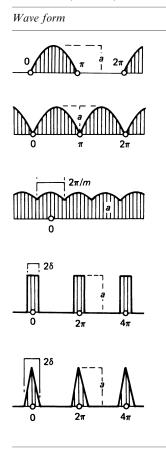
n	$J_n(1)$	$J_n(2)$	$J_n(3)$	$J_n(4)$	$J_n(5)$	$J_n(6)$	$J_n(7)$	$J_n(8)$	$J_n(9)$	<i>J<sub>n</sub></i> (10)	$J_n(11)$	<i>J<sub>n</sub></i> (12)	<i>J<sub>n</sub></i> (13)	<i>J<sub>n</sub></i> (14)	<i>J<sub>n</sub></i> (15)
0	0.7652	0.2239	-0.2601	-0.3971	-0.1776	0.1506	0.3001	0.1717	-0.0903	-0.2459	-0.1712	0.0477	0.2069	0.1711	-0.0142
1	0.4401	0.5767	0.3391	-0.0660	-0.3276	-0.2767	-0.0047	0.2346	0.2453	0.0435	-0.1768	-0.2234	-0.0703	0.1334	0.2051
2	0.1149	0.3528	0.4861	0.3641	0.0466	-0.2429	-0.3014	-0.1130	0.1448	0.2546	0.1390	-0.0849	-0.2177	-0.1520	0.0416
3	0.0196	0.1289	0.3091	0.4302	0.3648	0.1148	-0.1676	-0.2911	-0.1809	0.0584	0.2273	0.1951	0.0033	-0.1768	-0.1940
4	_	0.0340	0.1320	0.2811	0.3912	0.3567	0.1578	-0.1054	-0.2655	-0.2196	-0.0150	0.1825	0.2193	0.0762	-0.1192
5	_		0.0430	0.1321	0.2611	0.3621	0.3479	0.1858	-0.0550	-0.2341	-0.2383	-0.0735	0.1316	0.2204	0.1305
6	_		0.0114	0.0491	0.1310	0.2458	0.3392	0.3376	0.2043	-0.0145	-0.2016	-0.2437	-0.1180	0.0812	0.2061
7	_			0.0152	0.0534	0.1296	0.2336	0.3206	0.3275	0.2167	0.0184	-0.1703	-0.2406	-0.1508	0.0345
8	_				0.0184	0.0565	0.1280	0.2235	0.3051	0.3179	0.2250	0.0451	-0.1410	-0.2320	-0.1740
9	_					0.0212	0.0589	0.1263	0.2149	0.2919	0.3089	0.2304	0.0670	-0.1143	-0.2200
10	_			_		_	0.0235	0.0608	0.1247	0.2075	0.2804	0.3005	0.2338	0.0850	-0.0901
11	_			_		_		0.0256	0.0622	0.1231	0.2010	0.2704	0.2927	0.2357	0.0999
12	_								0.0274	0.0634	0.1216	0.1953	0.2615	0.2855	0.2367
13	_		—	_	_	_	_	_	0.0108	0.0290	0.0643	0.1201	0.1901	0.2536	0.2787
14	_		—	_	_	_	_	_	_	0.0119	0.0304	0.0650	0.1188	0.1855	0.2464
15	—		—	—	—	—	—	—	—	_	0.0130	0.0316	0.0656	0.1174	0.1813

Values below 0.01 not tabulated.

Table 1.10 Fourier series



#### Table 1.10 (continued)



Series

Rectified sine (half-wave): 
$$a\frac{1}{\pi} + a\frac{2}{\pi\varsigma} \left[ \frac{\pi \sin \theta_{\varsigma}}{4} - \frac{\cos 2\theta_{\varsigma}}{1 \cdot 3} - \frac{\cos 4\theta_{\varsigma}}{3 \cdot 5} - \frac{\cos 6\theta_{\varsigma}}{5 \cdot 7} - \cdots \right]$$

Rectified sine (full-wave): 
$$a\frac{2}{\pi} - a\frac{4}{\pi\varsigma} \left[ \frac{\cos 2\theta\varsigma}{1\cdot 3} + \frac{\cos 4\theta\varsigma}{3\cdot 5} + \frac{\cos 6\theta\varsigma}{5\cdot 7} + \frac{\cos 8\theta\varsigma}{7\cdot 9} + \cdots \right]$$

Rectified sine (*m*-phase): 
$$a\frac{m}{\pi\varsigma}\sin\frac{\pi\varsigma}{m} + a\frac{2m}{\pi\varsigma}\sin\frac{\pi\varsigma}{m}\left[\frac{\cos m\theta\varsigma}{m^2 - 1} - \frac{\cos 2m\theta\varsigma}{4m^2 - 1} + \frac{\cos 3m\theta\varsigma}{9m^2 - 1} - \cdots\right]$$

Rectangular pulse train: 
$$a\frac{\delta\varsigma}{\pi} + a\frac{2}{\pi\varsigma} \left[ \frac{\sin\delta\varsigma\cos\theta\varsigma}{1} + \frac{\sin2\delta\varsigma\cos2\theta\varsigma}{2} + \frac{\sin3\delta\varsigma\cos3\theta\varsigma}{3} + \cdots \right]$$
  
 $a\frac{\delta\varsigma}{\pi} + a\frac{2\delta\varsigma}{\pi\varsigma} \left[ \frac{\cos\theta\varsigma}{1} + \frac{\cos2\theta\varsigma}{2} + \frac{\cos3\theta\varsigma}{3} + \cdots \right]$  for  $\delta\varsigma \ll \langle \pi\varsigma \rangle$ 

Triangular pulse train: 
$$a\frac{\delta\varsigma}{2\pi\varsigma} + a\frac{4}{\pi\delta\varsigma} \left[ \frac{\sin^2(\frac{1}{2}\delta)}{1} \underbrace{\cos\theta}_{1} + \underbrace{\sin^2(\frac{1}{2}\delta)}_{4} \underbrace{\cos2\theta}_{2} + \underbrace{\sin^2(\frac{1}{2}\delta)}_{9} \underbrace{\cos3\theta}_{2} + \cdots \right] \leftarrow a\frac{\delta}{2\pi} + a\frac{\delta\varsigma}{\pi\varsigma} \underbrace{\cos\theta}_{1} + \cos2\theta + \cos3\theta + \cdots \right] \text{ for } \delta \ll \pi\varsigma$$

where the *a*s have the values either 1 or 0. Thus, if N = 19,  $19 = 16 + 2 + 1 = (2^4)1 + (2^3)0 + (2^2)0 + (2^1)1 + (2^0)1 = 10011$ in binary notation. The rules of addition and multiplication are

 $0 + 0 = 0, 0 + 1 = 1, 1 + 1 = 10; 0 \times 0 = 0, 0 \times 1 = 0, 1 \times 1 = 1$ 

#### 1.2.9 Power ratio

In communication networks the powers  $P_1$  and  $P_2$  at two specified points may differ widely as the result of amplification or attenuation. The power ratio  $P_1/P_2$  is more convenient in logarithmic terms.

*Neper* [Np] This is the natural logarithm of a voltage or current ratio, given by

$$a = 4n(V_1/V_2) \Leftarrow or \quad a = 4n(I_1/I_2) N_p$$

If the voltages are applied to, or the currents flow in, identical impedances, then the power ratio is

$$a = 4\mathbf{n}(V_1/V_2)^2 = 2 \ln(V_1/V_2) \Leftarrow$$

and similarly for current.

*Decibel* [dB] The power gain is given by the common logarithm  $lg(P_1/P_2)$  in bel [B], or most commonly by  $A = 10 \log(P_1/P_2)$  decibel [dB]. With again the proviso

that the powers are developed in identical impedances, the power gain is

$$A = 40 \log(P_1/P_2) = 40 \log(V_1/V_2)^2 = 20 \log(V_1/V_2) \text{ dB}$$

Table 1.13 gives the power ratio corresponding to a gain A (in dB) and the related identical-impedance voltage (or current) ratios. Approximately, 3 dB corresponds to a power ratio of 2, and 6 dB to a power ratio of 4. The decibel equivalent of 1 Np is 8.69 dB.

# 1.2.10 Matrices and vectors

# 1.2.10.1 Definitions

If  $a_{11}, a_{12}, a_{13}, a_{14}$ ... is a set of elements, then the rectangular array

$$\mathbf{A} = \left\{ \begin{bmatrix} a_{11} & a_{12} & a_{13} & a_{14} \dots a_{1n} \\ a_{21} & a_{22} & a_{23} & a_{24} \dots a_{2n} \\ a_{m1} & a_{m2} & a_{m3} & a_{m4} \dots a_{nm} \end{bmatrix} \right]$$

arranged in *m* rows and *n* columns is called an  $(m \times n)$  matrix. If m = n then **A** is *n*-square.

#### Table 1.11 Derivatives and integrals

d[f(x)]/dx	f(x)	$\int f(x) \cdot \mathbf{d}x$
1	x	$\frac{1}{2}x^2$
$nx^{n-1}$	$x^n (n \neq -1)$	$\frac{2}{x^{n+1}}/(n+1)$
$-1/x^{2}$	1/x	$\ln x$
1/x	$\ln x$	$x \ln x - x$
exp x	$\exp x$	$\exp x$
$\cos x$	sin x	$-\cos x$
$-\sin x$	$\cos x$	$\sin x$
$\sec^2 x$	tan x	$\ln(\sec x)$
$-\operatorname{cosec} x \cdot \operatorname{eot} x$	cosec x	$\ln(\tan\frac{1}{2}x)$
sec $x \cdot \tan x$	sec x	$\ln(\sec x + \tan x)$
$-\csc^2 x$	cot x	$\ln(\sin x)$
$1/\sqrt{(a^2-x^2)}$	$\arcsin(x/a)$	x $\arcsin(x/a) + \sqrt{a^2 - x^2}$
$-1/\sqrt{(a^2-x^2)}$	$\arccos(x/a)$	x $\arccos(x/a) - \sqrt{a^2 - x^2}$
$a/(a^2 + \langle x^2 \rangle)$	$\arctan(x/a)$	x $\arctan(x/a) - \frac{1}{2}a\ln(a^2 + x^2)$
$-a/x \sqrt{(x^2-a^2)}$	$\operatorname{arccosec}(x/a)$	$x \operatorname{arccosec}(x/a) + a \ln  x + \langle (x^2 - a^2) $
$a/x \sqrt{(x^2-a^2)}$	$\operatorname{arcsec}(x/a)$	x arcsec(x/a)-a ln   $x + \sqrt{(x^2 - a^2)}$
$-a/(a^2 + \langle \mathbf{x}^2 \rangle)$	$\operatorname{arccot}(x/a)$	$x \operatorname{arccot}(x/a) + \frac{1}{2}a\ln(a^2 + x^2)$
cosh x	sinh x	$\cosh x$
sinh x	cosh x	sinh x
$\operatorname{sech}^2 x$	tanh x	$\ln(\cosh x)$
$-\text{cosech } x \cdot \text{coth } x$	cosech x	$-\ln(\tanh\frac{1}{2}x)$
$-\operatorname{sech} x \cdot \operatorname{\mathfrak{sech}} x$	sech x	2 $\arctan(\exp x)$
$-\operatorname{cosech}^2 x$	coth x	$\ln(\sinh x)$
$1/\sqrt{(x^2+1)}$	arsinh x	x arsinh $x - \sqrt{1 + \alpha^2}$
$1/\sqrt{(x^2-1)}$	arcosh x	x arcosh $x - \sqrt{x^2 - 1}$
$1/(1-x^2)$	artanh x	x artanh $x + \frac{1}{2} \ln(1 - \frac{1}{2})$
$-1/x \sqrt{(x^2 + 4)}$	arcosech x	x arcosech $x + \operatorname{arsinh} x$
$-1/x \sqrt{(1-x^2)}$	arsech x	x arsech $x + \arcsin x$
$1/(1-x^2)$	arcoth x	x arcoth $x + \frac{1}{2} \ln(x^2 - 4)$
$u\frac{\mathrm{d}v}{\mathrm{d}x} + \frac{\mathrm{d}u}{\mathrm{d}x}$	$u(x) \cdot \#(x)$	$uv - \oint v \frac{\mathrm{d}u}{\mathrm{d}v} \mathrm{d}v$
1 du u dv	$\frac{u(x)}{v(x)}$	`]
$\frac{1}{v}\frac{dx}{dx} - \frac{dx}{v^2}\frac{dx}{dx}$	$\overline{v(x)}$	
$r \exp(xa) \times \sin(\omega x + \phi \varsigma + \langle \theta \rangle)$	$\exp(ax) \times \sin(\omega x + \nleftrightarrow)$	$(1/r)\exp(ax)\sin(\omega x + \phi - \theta)$ $r = \not\Leftarrow (\omega^2 + \not a^2) \ \theta \not \models \operatorname{arctan} (\omega/a)$

An ordered set of elements  $\mathbf{x} = [x_1, x_2, x_3 \dots x_n]$  is called an *n*-vector.

An  $(n \times 4)$  matrix is called a *column vector* and a  $(1 \times 4)$  matrix a *row vector*.

# 1.2.10.2 Basic operations

If  $\mathbf{A} = (a_{rs}), \mathbf{B} = (b_{rs}),$ 

- (i) Sum  $\mathbf{C} = \mathbf{A} + \mathbf{B}$  is defined by  $c_{rs} = \mathbf{a}_{\overline{r}s} + b_{rs}$ , for  $r = \mathbf{a}_{\overline{r}s} \dots m$ ;  $s = \mathbf{a}_{\overline{r}s} \dots n$ .
- (ii) Product If **A** is an  $(m \times q)$  matrix and **B** is a  $(q \times q)$  matrix, then the product  $\mathbf{C} = \mathbf{AB}$  is an  $(m \times q)$  matrix defined by  $(c_{rs}) = \sum_{p} a_{rp} b_{ps}$ ,  $p = 4 \dots q$ ;  $r = 4 \dots m$ ;  $s = 4 \dots n$ . If  $\mathbf{AB} = \mathbf{BA}$  then **A** and **B** are said to commute.
- (iii) Matrix-vector product If  $\mathbf{x} = \{\mathbf{x}_1 \dots \mathbf{x}_n\}$ , then  $\mathbf{b} = \mathbf{A}\mathbf{x}$  is defined by  $(b_r) = \sum_p a_{rp} x_p, p = \mathbf{A} \dots n; r = \mathbf{A} \dots m$ .
- (iv) Multiplication of a matrix by a (scalar) element If k is an element then  $\mathbf{C} = \mathbf{k}\mathbf{A} = \mathbf{k}k$  is defined by  $(c_{rs}) = \mathbf{k}(a_{rs})$ .
- (v) Equality If  $\mathbf{A} = \mathbf{B}$ , then  $(a_{ij}) = (\mathbf{b}_{ij})$ , for  $i = \mathbf{4} \dots n$ ;  $j = \mathbf{4} \dots m$ .

# 1.2.10.3 Rules of operation

- (i) Associativity  $\mathbf{A} + \langle \mathbf{B} + \langle \mathbf{C} \rangle = \langle \mathbf{A} + \langle \mathbf{B} \rangle + \langle \mathbf{C}, \mathbf{A}(\mathbf{BC}) = \langle \mathbf{A}\mathbf{B}\mathbf{C} = \mathbf{A}\mathbf{B}\mathbf{C}.$
- (ii) Distributivity A(B+C) = AB + AC, (B+C)A = BA + CA.
- (iii) *Identity* If U is the  $(n \times \#)$  matrix  $(\delta ij)$ ,  $i, j = \# \dots n$ , where  $\delta ij = \#$  if i = # and 0 otherwise, then U is the diagonal *unit matrix* and A U = A.
- (iv) *Inverse* If the product  $\mathbf{U} = \mathbf{A}\mathbf{B}$  exists, then  $\mathbf{B} = \mathbf{A}^{-1}$ , the inverse matrix of  $\mathbf{A}$ . If both inverses  $\mathbf{A}^{-1}$  and  $\mathbf{B}^{-1}$  exist, then  $(\mathbf{A} \mathbf{B})^{-1} = \mathbf{B}^{-1} \mathbf{A}^{-1}$ . (v) *Transposition* The transpose of  $\mathbf{A}$  is written as
- (v) Transposition The transpose of A is written as  $A^{T}$  and is the matrix whose rows are the columns of A. If the product C = AB exists then  $C^{T} = (AB)^{T} = (B^{T}A^{T})$ .
- (vi) Conjugate For  $\mathbf{A} = \{\mathbf{a}_{rs}\}$ , the congugate of  $\mathbf{A}$  is denoted by  $\mathbf{A}^* = \{\mathbf{a}_{rs}^*\}$ .
- (vii) Orthogonality Matrix A is orthogonal if  $AA^{T} = \mathcal{U}$ .

# Table 1.12 Laplace transforms

Definition	$f(t)$ from $t = \Theta +$	$\mathbf{F}(\mathbf{s}) = \mathbf{F}[f(t)] = \mathbf{F}[f(t) \cdot \mathbf{exp}(-\mathbf{s}t) \cdot $	
Sum	$af_1(t) + bf_2(t)$	$aF_1(s) + bF_2(s)$	
First derivative	(d/dt)f(t)	sF(s)-f(0-)	
<i>n</i> th derivative	$(\mathrm{d}^n/\mathrm{d}t^n)f(t)$	$s^{n}F(s)-s^{n-1}f(0-)-s^{n-2}f^{(1)}(0-)-$	$r^{(n-1)}(0-)$
Definite integral	$\int_{0-}^{T} f(t) \cdot \mathbf{d}t$	$\frac{1}{s}$ F(s)	
Shift by T	f(t-T)	$\exp(-sT) \cdot \mathbf{E}(s)$	
Periodic function (period $T$ )	f(t)	$\frac{1}{1 - \exp(-sT)} \int_{-\epsilon}^{T} \exp(-sT) \cdot f(t) \cdot dt$	
Initial value	$f(t), t \rightarrow 0 +$	$sF(s), s \rightarrow \infty \Leftarrow$	
Final value	$f(t), t \rightarrow \infty \Leftarrow$	$sF(s), s \rightarrow 0$	
Description	f(t)	F(s)	f(t) to base t
1. Unit impulse	$\delta(t)$	1	Area = 1
2. Unit step	H(t)	$\frac{1}{s}$	
3. Delayed step	H(t-T)	$\frac{\exp\left(-st\right)}{s}$	<b>→</b>
4. Rectangular pulse (duration $T$ )	H(t)-H(t-T)	$\frac{1 - \exp\left(-sT\right)}{s}$	
5. Unit ramp	t	$\frac{1}{s^2}$	
6. Delayed ramp	$(t-T)\mathbf{H}(t-T)$	$\frac{\exp\left(-sT\right)}{s^2}$	
7. <i>n</i> th-order ramp	ť <sup>n</sup>	$\frac{n!}{\mathbf{s}^{n+1}}$	
8. Exponential decay	$\exp(-\alpha t)$	$\frac{1}{s+\alpha\varsigma}$	
9. Exponential rise	$1 - \exp(-\alpha t)$	$\frac{\alpha\varsigma}{s(s+\alpha)}$	
10. Exponential ×≠	$t \exp(-\alpha t)$	$\frac{1}{(s+\alpha)^2}$	
11. Exponential $\times \not \in$	$t^n \exp(-\alpha t)$	$\frac{n!}{(s+\alpha)^{n+1}}$	
12. Difference of exponentials	$\exp(-\alpha t) - \exp(-\beta t)$	$\frac{\beta_{\varsigma-4\epsilon\varsigma}}{(s+4\epsilon)(s+4\epsilon)}$	

# 1/16 Units, mathematics and physical quantities

# Table 1.12 (continued)

Definition	f(t) from $t = 0 +$	$\mathbf{F}(\mathbf{s}) = L[f(t)] = \int_{0-}^{\infty \leftarrow} f(t) \cdot \exp(-\mathbf{s}t) \cdot dt$	
13. Sinusoidal	$\sin \omega t$	$\frac{\omega\varsigma}{s^2+\omega^2}$	
14. Phase-advanced sine	$\sin(\omega t + \phi)$	$\frac{\omega\cos\phi\varsigma+s\sin\phi\varsigma}{s^2+\omega^2}$	
15. Sine $\times t$	$t\sin\omega t$	$\frac{2\omega\mathrm{s}}{(\mathrm{s}^2+\omega^2)^2}$	
16. Exponentially decaying sine	$\exp(-\alpha t) \sin \omega t$	$\frac{\omega\varsigma}{\left(\mathbf{s}+\alpha\right)^2+\omega^2}$	
17. Cosinusoidal	$\cos \omega t$	$\frac{s}{s^2 + \omega^2}$	
18. Phase-advanced cosine	$\cos(\omega t + \phi)$	$\frac{s\cos\phi\varsigma-\omega\sin\phi\varsigma}{s^2+\omega^2}$	
19. Offset cosine	$1-\cos \omega t$	$\frac{\omega^2}{\mathrm{s}(\mathrm{s}^2+\omega^2)}$	
20. Cosine $\times t$	$t \cos \omega t$	$\frac{s^2-\omega^2}{(s^2+\omega^2)^2}$	
21. Exponentially decaying cosine	$\exp(-\alpha t)\cos\omega t$	$\frac{(s+\alpha)}{(s+\alpha)^2+\omega^2}$	
22. Trigonometrical function $G(t)$	$\sin \omega t - \omega t \cos \omega t$	$\frac{2\omega^3}{\left(s^2+\omega^2\right)^2}$	
23. Exponentially decaying trigonometrical function	$\exp(-\alpha t) \cdot \mathbf{G}(t)$	$\frac{2\omega^3}{\left[(s+\alpha)^2+\omega^2\right]^2}$	
24. Hyperbolic sine	$\sinh \omega t$	$\frac{\omega\varsigma}{\mathrm{s}^2-\omega^2}$	
25. Hyperbolic cosine	$\cosh \omega t$	$\frac{s}{s^2 - \omega^2}$	
26. Rectangular wave (period $T$ )	f(t)	$\frac{1 + \tanh(sT/4)}{2s}$	
27. Half-wave rectified sine $(T = 2\pi/\omega)$	f(t)	$\frac{\omega \exp{(sT/2)}\text{cosech}(sT/2)}{2(s^2 + \omega^2)}$	
28. Full-wave rectified sine $(T=2\pi/\omega)$	f(t)	$\frac{\omega \coth(sT/2)}{s^2 + \omega^2}$	

Table 1.13 Decibel gain: power and voltage ratios

A	$P_1/P_2$	$V_{1}/V_{2}$	A	$P_{1}/P_{2}$	$V_{1}/V_{2}$
0	1.000	1.000	9	7.943	2.818
0.1	1.023	1.012	10	10.00	3.162
0.2	1.047	1.023	12	15.85	3.981
0.3	1.072	1.032	14	25.12	5.012
0.4	1.096	1.047	16	39.81	6.310
0.6	1.148	1.072	18	63.10	7.943
0.8	1.202	1.096	20	100.0	10.00
1.0	1.259	1.122	25	316.2	17.78
1.2	1.318	1.148	30	1000	31.62
1.5	1.413	1.189	35	3162	56.23
2.0	1.585	1.259	40	$1.0 \times 10^{4}$	100.0
2.5	1.778	1.333	45	$3.2 \times 10^{4}$	177.8
3.0	1.995	1.413	50	$1.0 \times 10^{5}$	316.2
3.5	2.239	1.496	55	$3.2 \times 10^{5}$	562.3
4.0	2.512	1.585	60	$1.0 \times 10^{6}$	1 000
4.5	2.818	1.679	65	$3.2 \times 10^{6}$	1778
5.0	3.162	1.778	70	$1.0 \times 10^{7}$	3 1 6 0
6.0	3.981	1.995	80	$1.0 \times 10^{8}$	10 000
7.0	5.012	2.239	90	$1.0 \times 10^{9}$	31 620
8.0	6.310	2.512	100	$1.0  imes 10^{10}$	100 000

# 1.2.10.4 Determinant and trace

- (i) The *determinant* of a square matrix **A** denoted by  $|\mathbf{A}|$ , also det(**A**), is defined by the recursive formula  $|\mathbf{A}| = a_{11} \mathbf{M}_{11} a_{12} \mathbf{M}_{12} + a_{13} \mathbf{M}_{13} \cdots (\Leftarrow 1)^n a_{1n} \mathbf{M}_{1n}$  where  $\mathbf{M}_{11}$  is the determinant of the matrix with row 1 and column 1 missing,  $\mathbf{M}_{12}$  is the determinant of the matrix with row 1 and column 2 missing etc.
- (ii) The *Trace* of **A** is denoted by  $tr(\mathbf{A}) = \sum_i a_{ii}, i = 1, 2 \dots n$ .
- (iii) Singularity The square matrix  $\mathbf{A}$  is singular if det ( $\mathbf{A}$ ) = 0.
- (iv) The Characteristic Polynomial  $P(\lambda) = det(A \lambda U)$ .

# 1.2.10.5 Eigensystems

- *Eigenvalues* The eigenvalues of a matrix λ(A) are the *n* complex roots λ<sub>1</sub>(A), λ<sub>2</sub>(A)...λ<sub>n</sub>(A) of the characteristic polynomial det(A λU) = 0. Normally in most engineering systems there are no equal roots so the eigenvalues are distinct.
- (ii) Eigenvectors For any distinct eigenvalue λ<sub>i</sub> (A), there is an associated non-zero right eigenvector X<sub>i</sub> satisfying the homogeneous equations (A λ<sub>i</sub>U) X<sub>i</sub>=0, i=1,2...n. The matrix (A λ<sub>i</sub>U) is singular, however, because the det (A λ<sub>i</sub>U)=0; hence X<sub>i</sub> is not unique. In each set of equations (A λ<sub>i</sub>U) X<sub>i</sub>=0 one equation is redundant and only the relative values of the elements of X<sub>i</sub> can be determined. Thus the eigenvectors can be scaled arbitrarily, one element being assigned a value and the other elements determined accordingly from the remaining non-homogeneous equations.

The equations can be written also as  $\mathbf{A}\mathbf{X}_i = \lambda_i \mathbf{X}_i$ , or combining all eigenvalues and right eigenvectors,  $\mathbf{A}\mathbf{X} = \Lambda \mathbf{X}$ , where  $\Lambda$  is a diagonal matrix of the eigenvalues and  $\mathbf{X}$  is a square matrix containing all the right eigenvectors in corresponding order.

Since the eigenvalues of  $\mathbf{A}$  and  $\mathbf{A}^{T}$  are identical, for every eigenvalue  $\lambda_i$  associated with an eigenvector  $\mathbf{X}_i$  of  $\mathbf{A}$  there is also an eigenvector  $\mathbf{P}_i$  of  $\mathbf{A}^{T}$  such that  $\mathbf{A}^{T}\mathbf{P}_i = \lambda_i \mathbf{P}_i$ . Alternatively the eigenvector  $\mathbf{P}_i$  can be considered to be the *left eigenvector* of  $\mathbf{A}$  by transposing the equation to give  $\mathbf{P}_i^{T}\mathbf{A} = \lambda_i \mathbf{P}_i^{T}$ , or combining into one matrix equation,  $\mathbf{P}^{T}\mathbf{A} = \mathbf{P}^{T}\Lambda$ . *Reciprocal eigenvectors* Post-multiplying this last equation by the right eigenvector matrix **X** gives  $\mathbf{P}^{T}\mathbf{A}\mathbf{X} = \mathbf{P}^{T}\Lambda\mathbf{X}$ , which summarises the *n* sets of equations  $\mathbf{P}_{i}^{T}\mathbf{A}\mathbf{X}_{i} = \mathbf{P}_{i}^{T}\lambda_{i}\mathbf{X}_{i} = \lambda_{i}\mathbf{P}_{i}^{T}\mathbf{X}_{i} = k_{i}\lambda_{i}$ , where  $k_{i}$  is a scalar formed from the  $(1 \times n)$  by  $(n \times 1)$  vector product  $\mathbf{P}_{i}^{T}\mathbf{X}_{i}$ . With both  $\mathbf{P}_{i}$  and  $\mathbf{X}_{i}$  being scaled arbitrarily, re-scaling the left eigenvectors such that  $\mathbf{W}_{i} = (1/k_{i}) \mathbf{P}_{i}$ , gives  $\mathbf{W}_{i}^{T}\mathbf{X}_{j} = \leftarrow$  $\delta i j = 1$ , if i = j, and = 0 otherwise. In matrix form  $\mathbf{W}^{T}\mathbf{X} = \mathbf{U}$ , the unit matrix. The re-scaled left eigenvectors  $\mathbf{W}_{i}^{T}$  are said to be the reciprocal eigenvectors corresponding to the right eigenvectors  $\mathbf{X}_{i}$ .

(iii) Eigenvalue sensitivity analysis The change in the numerical value of  $\lambda_i$  with a change in any matrix A element  $\delta a_{rs}$  is to a first approximation given by  $\delta \lambda_i = (w_r)_i (x_s)_i \, \delta a_{rs}$  where  $(w_r)_i$  is the *r*-th element of the reciprocal eignvector  $\mathbf{W}_i$  corresponding to  $\lambda_i$  and  $(x_s)_i$  is the *s*-th element of the associated right eigenvector  $\mathbf{X}_i$ . In more compact form the sensitivity coefficients  $\delta \lambda_i / \delta a_{rs}$  or condition numbers of all *n* eigenvalues with respect to all elements of matrix A are expressible by the 1-term dyads  $\mathbf{S}_i = \mathbf{W}_i \mathbf{X}_i^T$ ,  $i = 1 \dots n$ .

$$\mathbf{S}_{i} = \left\{ \begin{array}{ccccc} \delta\lambda_{i}/\delta a_{11} & \delta\lambda_{i}/\delta a_{12} & \dots & \delta\lambda_{i}/\delta a_{1n} \\ \delta\lambda_{i}/\delta a_{21} & \delta\lambda_{i}/\delta a_{22} & \dots & \delta\lambda_{i}/\delta a_{2n} \\ \dots & \dots & \dots & \dots \\ \vdots & \vdots & \vdots & \vdots \\ \delta\lambda_{i}/\delta a_{n1} & \delta\lambda_{i}/\delta a_{n2} & \dots & \delta\lambda_{i}/\delta a_{nn} \end{array} \right\}$$

The matrix  $S_i$  is known as *i*-th *eigenvalue sensitivity* matrix.

- (iv) *Matrix functions* Transposed eigenvalue sensitivity matrices appear also in the dyadic expansion of a matrix and in matrix functions, thus  $\mathbf{A} = \sum_i \lambda_i \mathbf{X}_i \mathbf{W}_i^{\mathsf{T}} = \sum_i \lambda_i \mathbf{S}_i^{\mathsf{T}}$ ,  $i = 1 \dots n$ . Likewise  $[\mathbf{A}]^2 > [\sum_i \lambda_i \mathbf{S}_i^{\mathsf{T}}]^2 = \sum_i \lambda_i^2 \mathbf{S}_i^{\mathsf{T}}$  or in general  $[\mathbf{A}]^p = \sum_i \lambda_i^p \mathbf{S}_i^{\mathsf{T}}$ ; thus, for example,  $[\mathbf{A}]^{-1} = \leftarrow \sum_i \lambda_i^{-1} \mathbf{S}_i^{\mathsf{T}}$ .
- 1.2.10.6 Norms
- (i) Vector norms A scalar measure of the magnitude of a vector X with elements x<sub>1</sub>, x<sub>2</sub>...x<sub>n</sub>, is provided by a norm, the general family of norms being defined by ||X|| = [Σ<sub>i</sub> |x<sub>i</sub>|<sup>p</sup>]<sup>1/p</sup>. The usual norms are found from the values of p. If p = 1, ||X|| is the sum of the magnitudes of the elements, p = 2, ||X|| is Euclidean norm or square root of the sum of the squares of the magnitudes of the elements, p = infinity, ||X|| is the *infinity norm* or magnitude of the largest element.
- (ii) Matrix norms Several norms for matrices have also been defined, for matrix A two being the Euclidean norm, ||A||<sub>E</sub> = [Σ<sub>r</sub>Σ<sub>s</sub>|a<sub>rs</sub>|<sup>2</sup>]<sup>1/2</sup>, r = 1, 2...m; s = 1, 2...n, and the absolute norm, ||A|| = max<sub>r,s</sub> | a<sub>rs</sub>|.

# 1.3 Physical quantities

Engineering processes involve energy associated with physical materials to convert, transport or radiate energy. As energy has several natural forms, and as materials differ profoundly in their physical characteristics, separate technologies have been devised around specific processes; and materials may have to be considered macroscopically in bulk, or in microstructure (molecular, atomic and subatomic) in accordance with the applications or processes concerned.

# 1.3.1 Energy

Like 'force' and 'time', energy is a unifying concept invented to systematise physical phenomena. It almost defies precise definition, but may be described, as an aid to an intuitive appreciation.

*Energy* is the capacity for 'action' or *work*.

Work is the measure of the change in energy state.

State is the measure of the energy condition of a system.

*System* is the ordered arrangement of related physical entities or processes, represented by a *model*.

*Mode* is a description or mathematical formulation of the system to determine its *behaviour*.

*Behaviour* describes (verbally or mathematically) the energy processes involved in changes of state. Energy *storage* occurs if the work done on a system is recoverable in its original form. Energy *conversion* takes place when related changes of state concern energy in a different form, the process sometimes being reversible. Energy *dissipation* is an irreversible conversion into heat. Energy *transmission* and *radiation* are forms of energy transport in which there is a finite propagation time.

In a physical system there is an identifiable energy input  $W_i$  and output  $W_o$ . The system itself may store energy  $W_s$  and dissipate energy W. The energy conservation principle states that

 $W_{\rm i} = W_{\rm s} + W + W_{\rm o}$ 

Comparable statements can be made for energy changes  $\Delta w$  and for energy rates (i.e. powers), giving

 $\Delta w_{\rm i} = \Delta w_{\rm s} + \Delta w + \Delta w_{\rm o}$  and  $p_{\rm i} = p_{\rm s} + p + p_{\rm o}$ 

# 1.3.1.1 Analogues

In some cases the *mathematical* formulation of a system model resembles that of a model in a completely different physical system: the two systems are then analogues. Consider linear and rotary displacements in a simple mechanical system with the conditions in an electric circuit, with the following nomenclature: A mechanical element (such as a spring) of compliance k (which describes the displacement per unit force and is the inverse of the stiffness) has a displacement l = kf when a force f is applied. At a final force  $f_1$  the potential energy stored is  $W = \frac{1}{2}kf_1^2$ . For the rotary case,  $\theta = kM$  and  $W = \frac{1}{2}kM_1^2$ . In the electric circuit with a pure capacitance C, to which a p.d. v is applied, the charge is q = Cv and the electric energy stored at  $v_1$  is  $W = \frac{1}{2}Cv_1^2$ .

Use is made of these correspondences in mechanical problems (e.g. of vibration) when the parameters can be considered to be 'lumped'. An ideal transformer, in which the primary m.m.f. in ampere-turns  $i_1N_1$  is equal to the secondary m.m.f.  $i_2N_2$  has as analogue the simple lever, in which a force  $f_1$  at a point distant  $l_1$  from the fulcrum corresponds to  $f_2$  at  $l_2$  such that  $f_1l_1 = f_2l_2$ .

A simple series circuit is described by the equation v = L(di/dt) + Ri + q/C or, with *i* written as dq/dt,

$$v = L(\mathrm{d}^2 q/\mathrm{d}t^2) + R(\mathrm{d}q/\mathrm{d}t) + (1/C)q$$

A corresponding mechanical system of mass, compliance and viscous friction (proportional to velocity) in which for a displacement *l* the inertial force is m(du/dt), the compliance force is l/k and the friction force is ru, has a total force

$$f = m(d^2 l/dt^2) + r(dl/dt) + (1/k)l$$

Thus the two systems are expressed in identical mathematical form.

# 1.3.1.2 Fields

Several physical problems are concerned with 'fields' having stream-line properties. The eddyless flow of a liquid, the current in a conducting medium, the flow of heat from a high- to a low-temperature region, are fields in which representative lines can be drawn to indicate at any point the direction of

f m r k l	force [N] mass [kg] friction [N s/m] compliance [m/N] displacement [m] velocity [m/s]	$M \\ J \\ r \\ k \\ \theta_{\varsigma} \\ \omega_{\varsigma}$	torque [N m] inertia [kg m <sup>2</sup> ] friction [N m s/rad] compliance [rad/N m] displacement [rad] angular velocity [rad/s]	v L R C q i	voltage [V] inductance [H] resistance [Ω] capacitance [F] charge [C] current [A]
и	velocity [m/s]	$\omega \varsigma$	angular velocity [rad/s]	ĺ	current [A]

The force necessary to maintain a uniform linear velocity u against a viscous frictional resistance r is f = ur; the power is  $p = fu = u^2 r$  and the energy expended over a distance l is  $W = fut = u^2 rt$ , since l = ut. These are, respectively, the analogues of v = iR,  $p = vi = i^2 R$  and  $W = vit = i^2 Rt$  for the corresponding electrical system. For a constant angular velocity in a rotary mechanical system,  $M = \omega r$ ,  $p = M\omega = \omega^2 r$  and  $W = \omega^2 rt$ , since  $\theta = \omega t$ .

If a mass is given an acceleration du/dt, the force required is f = m(du/dt) and the stored kinetic energy at velocity  $u_1$ is  $W = \frac{1}{2}mu_1^2$ . For rotary acceleration,  $M = J(d\omega/dt)$  and  $W = \frac{1}{2}J\omega_1^2$ . Analogously the application of a voltage v to a pure inductor L produces an increase of current at the rate di/dt such that v = L(di/dt) and the magnetic energy stored at current  $i_1$  is  $W = \frac{1}{2}Li^2$ . the flow there. Other lines, orthogonal to the flow lines, connect points in the field having equal potential. Along these equipotential lines there is no tendency for flow to take place.

Static electric fields between charged conductors (having equipotential surfaces) are of interest in problems of insulation stressing. Magnetic fields, which in air-gaps may be assumed to cross between high-permeability ferromagnetic surfaces that are substantially equipotentials, may be studied in the course of investigations into flux distribution in machines. All the fields mentioned above satisfy Laplacian equations of the form

$$(\partial^2 V/\partial x^2) + (\partial^2 V/\partial y^2) + (\partial^2 V/\partial z^2) = 0$$

The solution for a physical field of given geometry will apply to other Laplacian fields of similar geometry, e.g.

System	Potential	Flux	Medium
current flow	voltage V	current $I$	conductivity $\sigma_{\varsigma}$
heat flow	temperature $\theta_{\varsigma}$	heat $q$	thermal conductivity $\lambda_{\varsigma}$
electric field	voltage V	electric flux $Q$	permittivity $\varepsilon_{\varsigma}$
magnetic field	m.m.f. F	magnetic flux $\phi_{\varsigma}$	permeability $\mu_{\varsigma}$

The ratio I/V for the first system would give the effective conductance G; correspondingly for the other systems,  $q/\theta\varsigma$ gives the thermal conductance, Q/V gives the capacitance and  $\Phi/F$  gives the permeance, so that if measurements are made in one system the results are applicable to all the others.

It is usual to treat problems as two-dimensional where possible. Several field-mapping techniques have been devised, generally electrical because of the greater convenience and problems, conductive methods include high-resistivity paper sheers, square-mesh 'nets' of resistors and electrolytic tanks. The tank is especially adaptable to three-dimensional cases of axial symmetry.

In the electrolytic *tank* a weak electrolyte, such as ordinary tap-water, provides the conducting medium. A scale model of the electrode system is set into the liquid. A low-voltage supply at some frequency between 50 Hz and 1 kHz is connected to the electrodes so that current flows through the electrolyte between them. A probe, adjustable in the horizontal plane and with its tip dipping vertically into the electrolyte, enables the potential field to be plotted. Electrode models are constructed from some suitable insulant (wood, paraffin wax, Bakelite, etc.), the electrode outlines being defined by a highly conductive material such as brass or copper. The metal is silver-plated to improve conductivity and reduce polarisation. Three-dimensional cases with axial symmetry are simulated by tilting the tank and using the surface of the electrolyte as a radial plane of the system.

The conducting-*sheet* analogue substitutes a sheet of resistive material (usually 'teledeltos' paper with silverpainted electrodes) for the electrolyte. The method is not readily adaptable to three-dimensional plots, but is quick and inexpensive in time and material.

The *mesh* or resistor-net analogue replaces a conductive continuum by a square mesh of equal resistors, the potential measurements being made at the nodes. Where the boundaries are simple, and where the 'grain size' is sufficiently small, good results are obtained. As there are no polarisation troubles, direct voltage supply can be used. If the resistors are made adjustable, the net can be adapted to cases of inhomogeneity, as when plotting a magnetic field in which permeability is dependent on flux density. Three-dimensional plots are made by arranging plane meshes in layers; the nodes are now the junctions of six instead of four resistors.

A stretched elastic membrane, depressed or elevated in appropriate regions, will accommodate itself smoothly to the differences in level: the height of the membrane everywhere can be shown to be in conformity with a two-dimensional Laplace equation. Using a rubber sheet as a membrane, the path of electrons in an electric field between electrodes in a vacuum can be investigated by the analogous paths of rolling bearing-balls. Many other useful analogues have been devised, some for the rapid solution of mathematical processes.

Recently considerable development has been made in point-by-point computer solutions for the more complicated field patterns in three-dimensional space.

# 1.3.2 Structure of matter

Material substances, whether solid, liquid or gaseous, are conceived as composed of very large numbers of *molecules*. A molecule is the smallest portion of any substance which cannot be further subdivided without losing its characteristic material properties. In all states of matter molecules are in a state of rapid continuous motion. In a *solid* the molecules are relatively closely 'packed' and the molecules, although rapidly moving, maintain a fixed mean position. Attractive forces between molecules account for the tendency of the solid to retain its shape. In a *liquid* the molecules are less closely packed and there is a weaker cohesion between them, so that they can wander about with some freedom within the liquid, which consequently takes up the shape of the vessel in which it is contained. The molecules in a *gas* are still more mobile, and are relatively far apart. The cohesive force is very small, and the gas is enabled freely to contract and expand. The usual effect of heat is to increase the intensity and speed of molecular activity so that 'collisions' between molecules increase, so that the substance attempts to expand, producing internal pressure if the expansion is resisted.

Molecules are capable of further subdivision, but the resulting particles, called *atoms*, no longer have the same properties as the molecules from which they came. An atom is the smallest portion of matter than can enter into chemical combination or be chemically separated, but it cannot generally maintain a separate existence except in the few special cases where a single atom forms a molecule. A molecule may consist of one, two or more (sometimes many more) atoms of various kinds. A substance whose molecules are composed entirely of atoms of the same kind is called an *element*. Where atoms of two or more kinds are present, the molecule is that of a chemical *compound*. At present over 100 elements are recognised (*Table 1.14*: the atomic mass number A is relative to 1/12 of the mass of an element of carbon-12).

If the element symbols are arranged in a table in ascending order of atomic number, and in columns ('groups') and rows ('periods') with due regard to associated similarities, *Table 1.15* is obtained. Metallic elements are found on the left, non-metals on the right. Some of the correspondences that emerge are:

Group 1a: Alkali metals
(Li 3, Na 11, K 19, Rb 37, Cs 55, Fr 87)
2a: Alkaline earths
(Be 4, Mg 12, Ca 20, Sr 38, Ba 56, Ra 88)
1b: Copper group (Cu 29, Ag 47, Au 79)
6b: Chromium group (Cr 24, Mo 42, W 74)
7a: Halogens (F 9, Cl 17, Br 35, I 53, At 85)
0: Rare gases
(He 2, Ne 10, Ar 18, Kr 36, Xe 54, Rn 86)
3a–6a: Semiconductors
(B 5, Si 16, Ge 32, As 33, Sb 51, Te 52)

In some cases a horizontal relation obtains as in the transition series (Sc 21...Ni 28) and the heavy-atom rare earth and actinide series. The explanation lies in the structure of the atom.

#### 1.3.2.1 Atomic structure

The original Bohr model of the hydrogen atom was a central nucleus containing almost the whole mass of the atom, and a single *electron* orbiting around it. Electrons, as small particles of negative electric charge, were discovered at the end of the nineteenth century, bringing to light the complex structure of atoms. The hydrogen nucleus is a *proton*, a mass having a charge equal to that of an electron, but positive. Extended to all elements, each has a nucleus comprising mass particles, some (*protons*) with a positive charge, others (*neutrons*) with no charge. The atomic *mass number A* is the total number of protons and neutrons in the nucleus; the *atomic number Z* is the number of positive charges, and the normal number of orbital electrons. The nuclear structure is not known, and the forces that bind the protons against their mutual attraction are conjectural.

The hydrogen atom (Figure 1.4) has one proton (Z=4)and one electron in an orbit formerly called the K shell. Helium (Z=2) has two protons, the two electrons occupying the K shell which, by the Pauli exclusion principle, cannot have more than two. The next element in order is lithium (Z=3), the third electron in an outer L shell. With elements of increasing atomic number, the electrons are added to the L shell until it holds a maximum of 8, the surplus then occupying the M shell to a maximum of 18. The number of 'valence' electrons (those in the outermost shell) determines the physical and chemical properties of the element. Those with completed outer shells are 'stable'.

*Isotopes* An element is often found to be a mixture of atoms with the same chemical property but different atomic masses: these are isotopes. The isotopes of an element must have the same number of electrons and protons, but differ in the number of neutrons, accounting for the non-integral average mass numbers. For example, neon comprises 90.4% of mass number 20, with 0.6% of 21 and 9.0% of mass number 22, giving a resultant mass number of 20.18.

*Energy states* Atoms may be in various energy states. Thus, the filament of an incadescent lamp may emit light when excited by an electric current but not when the current is switched off. Heat energy is the kinetic energy of the atoms of a heated body. The more vigorous impact of atoms may not always shift the atom as a whole, but may shift an electron from one orbit to another of higher energy level within the atom. This position is not normally stable, and the electron gives up its momentarily acquired potential energy by falling back to its original level, releasing the energy as a light quantum or photon.

*Ionisation* Among the electrons of an atom, those of the outermost shell are unique in that, on account of all the electron charges on the shells between them and the nucleus, they are the most loosely bound and most easily *removable*. In a variety of ways it is possible so to excite an atom that one of the outer electrons is torn away, leaving the atom *ionised* or converted for the time into an *ion* with an effective positive charge due to the unbalanced electrical state it has acquired. Ionisation may occur due to impact by other fast-moving particles, by irradiation with rays of suitable wavelength and by the application of intense electric fields.

# 1.3.2.2 Wave mechanics

The fundamental laws of optics can be explained without regard to the nature of light as an electromagnetic wave phenomenon, and photoelectricity emphasises its nature as a stream or ray of corpuscles. The phenomena of diffraction or interference can only be explained on the wave concept. Wave mechanics correlates the two apparently conflicting ideas into a wider concept of 'waves of matter'. Electrons, atoms and even molecules participate in this duality, in that their effects appear sometimes as corpuscular, sometimes as of a wave nature. Streams of electrons behave in a corpuscular fashion in photoemission, but in certain circumstances show the diffraction effects familiar in wave action. Considerations of particle mechanics led de Broglie to write several theoretic papers (1922-1926) on the parallelism between the dynamics of a particle and geometrical optics, and suggested that it was necessary to admit that classical dynamics could not interpret phenomena involving energy quanta. Wave mechanics was established by Schrödinger in 1926 on de Broglie's conceptions.

When electrons interact with matter, they exhibit wave properties: in the free state they act like particles. Light has a similar duality, as already noted. The hypothesis of de Broglie is that a particle of mass m and velocity u has wave

 Table 1.14
 Elements (Z, atomic number; A, atomic mass;

 KLMNOPQ, electron shells)
 Image: Comparison of the shell of the s

Ζ	Name and symbol		Α		Shells
				K	L
1	Hydrogen	Η	1.008	1	
2	Helium	He	4.002	2	
3	Lithium	Li	6.94	2	1
4	Beryllium	Be	9.02	2	2
5	Boron	B	10.82	2	3
6	Carbon	C	12	2	4
7	Nitrogen	N	14.01	2	5
8 9	Oxygen Fluorine	O F	16.00	2	6 7
10	Neon	г Ne	$19.00 \\ 20.18$	2 2	8
10	INCOL	140	20.10		
11	Sodium	Na	22.99	KL 10	M N 1 —
12	Magnesium	Mg	24.32	10	2 —
13	Aluminium	Al	26.97	10	<u>3</u> —
14	Silicon	Si	28.06	10	4 —
15	Phosphorus	P	31.02	10	5 —
16	Sulphur	S	32.06	10	6 —
17	Chlorine	Cl	35.46	10	7 —
18	Argon	Ar	39.94	10	8 —
19	Potassium	Κ	39.09	10	8 1
20	Calcium	Ca	40.08	10	8 2
21	Scandium	Sc	45.10	10	9 2
22	Titanium	Ti	47.90	10	10 2
23	Vanadium	V	0.95	10	11 2
24	Chromium	Cr	52.01	10	13 1
25	Manganese	Mn	54.93	10	13 2
26	Iron	Fe	55.84	10	14 2 15 2
27 28	Cobalt Nickel	Co Ni	58.94 58.69	10 10	$\begin{array}{ccc} 15 & 2 \\ 16 & 2 \end{array}$
28 29	Copper	Cu	63.57	10	18 1
30	Zinc	Zn	65.38	10	18 2
31	Gallium	Ga	69.72	10	18 3
32	Germanium	Ge	72.60	10	18 4
33	Arsenic	Ās	74.91	10	18 5
34	Selenium	Se	78.96	10	18 6
35	Bromine	Br	79.91	10	18 7
36	Krypton	Kr	83.70	10	18 8
				KLM	ΝΟ
37	Rubidium	Rb	85.44	28	8 1
38	Strontium	Sr	87.63	28	8 2
39	Yttrium	Y	88.92	28	$\begin{array}{rrrr} 8 & 2 \\ 9 & 2 \\ 10 & 2 \end{array}$
40	Zirconium	Zr	91.22	28	
41	Niobium	Nb	92.91	28	12 1 13 1
42 43	Molybdenum	Mo	96.0	28 28	
43 44	Technetium Ruthenium	Tc Ru	99.0 101.7	28 28	14 1 15 1
45	Rhodium	Rh	101.7	28	16 1
46	Palladium	Pd	102.9	28	18 —
47	Silver	Ag	107.9	28	18 1
48	Cadmium	Cď	112.4	28	18 2
49	Indium	In	114.8	28	18 3
50	Tin	Sn	118.7	28	18 4
51	Antimony	Sb	121.8	28	18 5
52	Tellurium	Te	127.6	28	18 6
53	Iodine	Ι	126.9	28	18 7
54	Xenon	Xe	131.3	28	18 8
				KLM	NOD
				NLIV	NUP
55	Caesium	Cs	132.9	28	N O P 18 8 1
55 56	Caesium Barium	Cs Ba	132.9 137.4		

Table 1.14 (continued)

Ζ	Name and symbol		A			Sh	ells
				KLM	N	0	Р
57	Lanthanum	La	138.9	28	18	9	2
58	Cerium	Ce	140.1	28	19	9	2 2 2
59	Praseodymium	Pr	140.9	28	21	8	2
60	Neodymium	Nd	144.3	28	22	8	2
61	Promethium	Pm	147.0	28	23	8	2
62	Samarium	Sm	150.4	28	24	8	2
63	Europium	Eu	152.0	28	25	8	2
64	Gadolinium	Gd	157.3	28	25	9	2
65	Terbium	Tb	159.2	28	27	8	2
66	Dysprosium	Dy	162.5	28	28	8	2
67	Holmium	Ho	163.5	28	29	8	2
68	Erbium	Er	167.6	28	30	8	2
69	Thulium	Tm	169.4	28	31	8	$\overline{2}$
70	Ytterbium	Yb	173.0	28	32	8	$\overline{2}$
71	Lutecium	Lu	175.0	28	32	9	2
72	Hafnium	Hf	178.6	28	32	10	$     \begin{array}{c}       2 \\     $
73	Tantalum	Ta	181.4	28	32	11	2
74	Tungsten	W	184.0	28	32	12	2 2 2
75	Rhenium	Re	186.3	28	32	13	2
76	Osmium	Os	191.5	28	32	14	$\frac{2}{2}$
77	Iridium	Ir	191.5	28	32	15	$\frac{2}{2}$
78	Platinum	Pt	195.1	28	32	17	1
79	Gold	Au	195.2	28	32	18	1
80		Hg	200.6	28	32	18	2
81	Mercury Thallium	Tl	200.0	28	32	18	$\frac{2}{3}$
82	Lead	Pb	204.4	28 28	32	18	4
82 83		Bi			32 32		5
	Bismuth	Ро	209.0	28	32 32	18	5 6
84	Polonium		210.0	28		18	
85	Astatine	At	211.0	28	32	18	7
86	Radon	Rn	222.0	28	32	18	8
07	<b>.</b> .	Б	222.0	KLMN	0	P	Q
87	Francium	Fr	223.0	60	18	8	1
88	Radium	Ra	226.0	60	18	8	2 2 2
89	Actinium	Ac	227.0	60	18	9	2
90	Thorium	Th	232.0	60	18	10	2
91	Protoactinium	Pa	231.0	60	20	9	2
92	Uranium	U	238.0	60	21	9	2
93	Neptunium	Np	237.0	60	22	9	2
94	Plutonium	Pu	239.0	60	24	8	2
95	Americium	Am	243.0	60	25	8	2
96	Curium	Cm	247.0	60	25	9	
97	Berkelium	Bk	247.0	60	26	9	2
98	Californium	Cf	251.0	60	28	8	2
99	Einsteinium	Es	254.0	60	29	8	2
100	Fermium	Fm	257.0	60	30	8	2
101	Mendelevium	Md	257.0	60	31	8	2
102	Nobelium	No	254.0	60	32	8	2
				60	22	0	2
103	Lawrencium	Lr	256.0	60	32	9	2
103 104	Lawrencium Kurchatovium	Lr Ku	256.0	60	32	9	Ζ

properties with a wavelength  $\lambda = \frac{h}{mu}$ , where *h* is the Planck constant,  $h = \frac{6}{5.626} \times \frac{40}{34}$  J s. The mass *m* is relativistically affected by the velocity.

When electron waves are associated with an atom, only certain fixed-energy states are possible. The electron can be raised from one state to another if it is provided, by some external stimulus such as a photon, with the necessary energy difference  $\Delta w$  in the form of an electromagnetic wave of wavelength  $\lambda = 4c/\Delta w$ , where c is the velocity of free space radiation (3 × 40<sup>8</sup> m/s). Similarly, if an electron

falls from a state of higher to one of lower energy, it emits energy  $\Delta w$  as radiation. When electrons are raised in energy level, the atom is *excited*, but not ionised.

#### 1.3.2.3 Electrons in atoms

Consider the hydrogen atom. Its single electron is not located at a fixed point, but can be anywhere in a region near the nucleus with some probability. The particular region is a kind of shell or cloud, of radius depending on the electron's energy state.

With a nucleus of atomic number Z, the Z electrons can have several possible configurations. There is a certain radial pattern of electron probability cloud distribution (or shell pattern). Each electron state gives rise to a cloud pattern, characterised by a definite energy level, and described by the series of quantum numbers  $n, l, m_l$  and  $m_s$ . The number n(=4, 2, 3, ...) is a measure of the energy level; l(=0, 1, 2, ...) is concerned with angular momentum;  $m_l$  is a measure of the component of angular momentum in the direction of an applied magnetic field; and  $m_s$  arises from the electron spin. It is customary to condense the nomenclature so that electron states corresponding to  $l = \Phi$ . 1, 2 and 3 are described by the letters s, p, d and f and a numerical prefix gives the value of n. Thus boron has 2 electrons at level 1 with  $l = \Theta$ , two at level 2 with  $l = \Theta$ , and one at level 3 with l=4: this information is conveyed by the description  $(1s)^2(2s)^2(2p)^1$ .

The energy of an atom as a whole can vary according to the electron arrangement. The most stable state is that of minimum energy, and states of higher energy content are excited. By Pauli's exclusion principle the maximum possible number of electrons in states 1,  $\overline{2}$ , 3, 4, ..., n are 2, 8, 18,  $32, \ldots, 2n^2$ , respectively. Thus, only 2 electrons can occupy the 1s state (or K shell) and the remainder must, even for the normal minimum-energy condition, occupy other states. Hydrogen and helium, the first two elements, have, respectively, 1 and 2 electrons in the 1-quantum (K) shell; the next, lithium, has its third electron in the 2-quantum (L) shell. The passage from lithium to neon results in the filling up of this shell to its full complement of 8 electrons. During the process, the electrons first enter the 2s subgroup, then fill the 2p subgroup until it has 6 electrons, the maximum allowable by the exclusion principle (see *Table 1.14*).

Very briefly, the effect of the electron-shell filling is as follows. Elements in the same chemical family have the same number of electrons in the subshell that is incompletely filled. The rare gases (He, Ne, Ar, Kr, Xe) have no uncompleted shells. Alkali metals (e.g. Na) have shells containing a single electron. The alkaline earths have two electrons in uncompleted shells. The good conductors (Ag, Cu, Au) have a single electron in the uppermost quantum state. An irregularity in the ordered sequence of filling (which holds consistently from H to Ar) begins at potassium (K) and continues to Ni, becoming again regular with Cu, and beginning a new irregularity with Rb.

The electron of a hydrogen atom, normally at level 1, can be raised to level 2 by endowing it with a particular quantity of energy most readily expressed as 10.2 eV.  $(1 \text{ eV} = 4 \text{ electron-volt} = 4.6 \times 40^{-19} \text{ J}$  is the energy acquired by a free electron falling through a potential difference of 1 V, which accelerates it and gives it kinetic energy.) The 10.2 V is the *first excitation potential* for the hydrogen atom. If the electron is given an energy of 13.6 eV, it is freed from the atom, and 13.6 V is the *ionisation potential*. Other atoms have different potentials in accordance with their atomic arrangement.

Table 1.15	Elements:	periodic	table
------------	-----------	----------	-------

Periods									Grou	ps								
	1a	2a	3b	4b	5b	6b	7b	8b	8b	8b	1b	2b	3a	4a	5a	6a	7a	0
Ι	1 H				Meta	ls								Non-1	netals			2 He
II	3 Li	4 Be											5 B	6 C	7 N	8 O	9 F	10 Ne
III	11 Na	12 Mg			Trans	sitions							13 Al	14 Si	15 P	16 S	17 Cl	18 Ar
IV	19 K	20 Ca	21 Sc	22 Ti	23 V	24 Cr	25 Mn	26 Fe	27 Co	28 Ni	29 Cu	30 Zn	31 Ga	32 Ge	33 As	34 Se	35 Br	36 Kr
v	37 Rb	38 Sr	39 Y	40 Zr	41 Nb	42 Mo	43 Tc	44 Ru	45 Rh	46 Pd	47 Ag	48 Cd	49 In	50 Sn	51 Sb	52 Te	53 I	54 Xe
VI	55 Cs	56 Ba	57 La	72 Hf	73 Ta	74 W	75 Re	76 Os	77 Ir	78 Pt	79 Au	80 Hg	81 Tl	82 Pb	83 Bi	84 Po	85 At	86 Rn
VII	87 Fr	88 Ra	89 Ac															
		Rare e	earths		58 Ce	59 Pr	60 Nd	61 Pm	62 Sm	63 Eu	64 Gd	65 Tb	66 Dy	67 Ho	68 Er	69 Tm	70 Yb	71 Lu
		Actini	des		90 Th	91 Pa	92 U	93 Np	94 Pu	95 Am	96 Cm	97 Bk	98 Cf	99 Es	100 Fm	101 Md	102 No	103 Lr

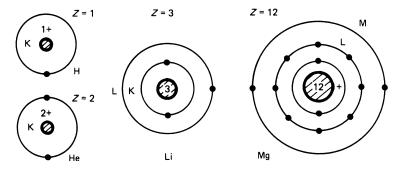


Figure 1.4 Atomic structure

# 1.3.2.4 Electrons in metals

An approximation to the behaviour of metals assumes that the atoms lose their valency electrons, which are free to wander in the ionic lattice of the material to form what is called an electron gas. The sharp energy levels of the free atom are broadened into wide bands by the proximity of others. The potential within the metal is assumed to be smoothed out, and there is a sharp rise of potential at the surface which prevents the electrons from escaping: there is a potential-energy step at the surface which the electrons cannot normally overcome: it is of the order of 10 eV. If this is called W, then the energy of an electron wandering within the metals is  $-W + \frac{1}{2}mu^2$ .

The electrons are regarded as undergoing continual collisions on account of the thermal vibrations of the lattice, and on Fermi-Dirac statistical theory it is justifiable to treat the energy states (which are in accordance with Pauli's principle) as forming an energy continuum. At very low temperatures the ordinary classical theory would suggest that electron energies spread over an almost zero range, but the exclusion principle makes this impossible and even at absolute zero of temperature the energies form a continuum, and physical properties will depend on how the electrons are distributed over the upper levels of this energy range.

*Conductivity* The interaction of free electrons with the thermal vibrations of the ionic lattice (called 'collisions' for brevity) causes them to 'rebound' with a velocity of random direction but small compared with their average velocities as particles of an electron gas. Just as a difference of electric potential causes a drift in the general motion, so a difference of temperature between two parts of a metal carries energy from the hot region to the cold, accounting for thermal conduction and for its association with electrical conductivity. The free electron theory, however, is inadequate to explain the dependence of conductivity on crystal axes in the metal.

At absolute zero of temperature (zero  $K = -273^{\circ}C$ ) the atoms cease to vibrate, and free electrons can pass through the lattice with little hindrance. At temperatures over the range 0.3–10K (and usually round about 5 K) the resistance of certain metals, e.g. Zn, Al, Sn, Hg and Cu, becomes substantially zero. This phenomenon, known as *superconductivity*, has not been satisfactorily explained.

Superconductivity is destroyed by moderate magnetic fields. It can also be destroyed if the current is large enough to produce at the surface the same critical value of magnetic field. It follows that during the superconductivity phase the current must be almost purely superficial, with a depth of penetration of the order of  $10 \,\mu\text{m}$ .

#### Table 1.16 Physical properties of metals

Approximate general properties at normal temperatures:

- $\delta\varsigma$  density [kg/m<sup>3</sup>]
- *E* elastic modulus [GPa]
- *e* linear expansivity  $[\mu m/(m K)]$
- c specific heat capacity [kJ/(kgK)]
- k thermal conductivity [W/(m K)]  $T_{\rm m}$  melting point [K]
- [K]  $\rho_{\varsigma}$  resistivity  $[n\Omega m]$

 $\alpha_{\varsigma}$  resistance-temperature coefficient [m $\Omega/(\Omega K)$ ]

Metal	$\delta\varsigma$	Ε	е	с	k	$T_{\rm m}$	$ ho\varsigma$	$\alpha\varsigma$
Pure metals								
4 Beryllium	1840	300	120	1700	170	1560	33	9.0
11 Sodium	970		71	710	130	370	47	5.5
12 Magnesium	1740	44	26	1020	170	920	46	3.8
13 Aluminium	2700	70	24	900	220	930	27	4.2
19 Potassium	860		83	750	130	340	67	5.4
20 Calcium	1550		22	650	96	1120	43	4.2
24 Chromium	7100	25	8.5	450	43	2170	130	3.0
26 Iron	7860	220	12	450	75	1810	105	6.5
27 Cobalt	8800	210	13	420	70	1770	65	6.2
28 Nickel	8900	200	13	450	70	1730	78	6.5
29 Copper	8930	120	16	390	390	1360	17	4.3
30 Zinc	7100	93	26	390	110	690	62	4.1
42 Molybdenum	10 200		5	260	140	2890	56	4.3
47 Silver	10 500	79	19	230	420	1230	16	3.9
48 Cadmium	8640	60	32	230	92	590	75	4.0
50 Tin	7300	55	27	230	65	500	115	4.3
73 Tantalum	16 600	190	6.5	140	54	3270	155	3.1
74 Tungsten	19 300	360	4	130	170	3650	55	4.9
78 Platinum	21 500	165	9	130	70	2050	106	3.9
79 Gold	19 300	80	14	130	300	1340	23	3.6
80 Mercury	13 550		180	140	10	230	960	0.9
82 Lead	11 300	15	29	130	35	600	210	4.1
83 Bismuth	9800	32	13	120	9	540	1190	4.3
92 Uranium	18 700	13	—	120		1410	220	2.1
Alloys								
Brass (60 Cu, 40 Zn)	8500	100	21	380	120	1170	60	2.0
Bronze (90 Cu, 10 Sn)	8900	100	19	380	46	1280		
Constantan	8900	110	15	410	22	1540	450	0.05
Invar (64 Fe, 36 Ni)	8100	145	2	500	16	1720	100	2.0
Iron, soft (0.2 C)	7600	220	12	460	60	1800	140	
Iron cast (3.5 C, 2.5 Si)	7300	100	12	460	60	1450		
Manganin	8500	130	16	410	22	1270	430	0.02
Steel (0.85 C)	7800	200	12	480	50	1630	180	

Electron emission A metal may be regarded as a potential 'well' of depth -V relative to its surface, so that an electron in the lowest energy state has (at absolute zero temperature) the energy  $W = 4\epsilon$  (of the order 10 eV): other electrons occupy levels up to a height  $\epsilon^*$  (5–8 eV) from the bottom of the 'well'. Before an electron can escape from the surface it must be endowed with an energy not less than  $\phi = 4W - \epsilon^*$ , called the *work function*.

Emission occurs by *surface irradiation* (e.g. with light) of frequency v if the energy quantum hv of the radiation is at least equal to  $\phi$ . The threshold of photoelectric emission is therefore with radiation at a frequency not less than  $v = \frac{\phi}{h}$ .

Emission takes place at *high temperatures* if, put simply, the kinetic energy of electrons normal to the surface is great enough to jump the potential step W. This leads to an expression for the emission current *i* in terms of temperature T, a constant A and the thermionic work function  $\phi$ :

$$i = \not A T^2 \exp(-\phi/kT) \not \in$$

Electron emission is also the result of the application of a *high electric field intensity* (of the order 1-10 GV/m) to a

metal surface; also when the surface is bombarded with electrons or ions of sufficient kinetic energy, giving the effect of *secondary* emission.

*Crystals* When atoms are brought together to form a crystal, their individual sharp and well-defined energy levels merge into energy *bands*. These bands may overlap, or there may be gaps in the energy levels available, depending on the lattice spacing and interatomic bonding. Conduction can take place only by electron migration into an empty or partly filled band; filled bands are not available. If an electron acquires a small amount of energy from the externally applied electric field, and can move into an available empty level, it can then contribute to the conduction process.

# 1.3.2.5 Insulators

In this case the 'distance' (or energy increase  $\Delta w$  in electronvolts) is too large for moderate electric applied fields to endow electrons with sufficient energy, so the material remains an insulator. High temperatures, however, may

#### 1/24 Units, mathematics and physical quantities

#### Table 1.17 Physical properties of non-metals

Approximate general properties:

- $\delta\varsigma$  density [kg/m<sup>3</sup>]
- *e* linear expansivity  $[\mu m/(m K)]$
- c specific heat capacity [kJ/(kgK)]
- k thermal conductivity [W/(m K)]

 $T_{\rm m}$  melting point [K]

 $\rho_{\varsigma}$  resistivity [M $\Omega$  m]

 $\epsilon_{\rm r}$  relative permittivity [-]

Material	$\delta\varsigma$	е	с	k	$T_{\rm m}$	$ ho\varsigma$	$\varepsilon_{\mathbf{F}}$
Asbestos (packed)	580		0.84	0.19			3
Bakelite	1300	30	0.92	0.20		0.1	7
Concrete (dry)	2000	10	0.92	1.70			
Diamond	3510	1.3	0.49	165	4000	$10^{7}$	
Glass	2500	8	0.84	0.93		$10^{6}$	8
Graphite	2250	2	0.69	160	3800	$10^{-11}$	
Marble	2700	12	0.88	3		$10^{3}$	8.5
Mica	2800	3	0.88	0.5		$10^{8}$	7
Nylon	1140	100	1.7	0.3	_		
Paper	900			0.18	_	$10^{4}$	2
Paraffin wax	890	110	2.9	0.26		$10^{9}$	2
Perspex	1200	80	1.5	1.9		$10^{14}$	3
Polythene	930	180	2.2	0.3			2.3
Porcelain	2400	3.5	0.8	1.0	1900	$10^{6}$	6
Quartz (fused)	2200	0.4	0.75	0.22	2000	$10^{14}$	3.8
Rubber	1250	_	1.5	0.15		107	3
Silicon	2300	7	0.75	_	1690	0.1	2.7

# Table 1.18 Physical properties of liquids

Average values at 20°C (293 K):

 $\delta \varsigma$  density [kg/m<sup>3</sup>]

v viscosity [mPa s]

*e* cubic expansivity  $[10^{-3}/K]$ 

 $T_{\rm m}$  melting point [K]

thermal conductivity [W/(mK)]

 $T_{\rm b}$  boiling point [K]

k

c specific heat capacity [kJ/(kg K)]

 $\varepsilon_{\mathbf{F}}$  relative permittivity [-]

Liquid		$\delta\varsigma$	v	е	с	k	$T_{\rm m}$	$T_{\rm b}$	$\varepsilon_{\mathbf{S}}$
Acetone	(CH <sub>3</sub> ) <sub>2</sub> CO	792	0.3	1.43	2.2	0.18	178	329	22
Benzine	$C_6H_6$	881	0.7	1.15	1.7	0.14	279	353	2.3
Carbon disulphide	$CS_2$	1260	0.4	1.22	1.0	0.14	161	319	2.6
Carbon tetrachloride	CCĨ4	1600	1.0	1.22	0.8	0.10	250	350	2.2
Ether	$(C_2H_5)_2O$	716	0.2	1.62	2.3	0.14	157	308	4.3
Glycerol	$C_3H_5(OH)_3$	1270	1500	0.50	2.4	0.28	291	563	56
Methanol	CH <sub>3</sub> OH	793	0.6	1.20	1.2	0.21	175	338	32
Oil	_	850	85	0.75	1.6	0.17			3.0
Sulphuric acid	$H_2SO_4$	1850	28	0.56	1.4		284	599	
Turpentine	$C_{10}H_{16}$	840	1.5	0.10	1.8	0.15	263	453	2.3
Water	$H_2O$	1000	1.0	0.18	4.2	0.60	273	373	81

result in sufficient thermal agitation to permit electrons to 'jump the gap'.

# attract an electron from a filled band to leave a 'hole', or electron deficiency, the movement of which corresponds to the movement of a positive charge (p-type).

# 1.3.2.6 Semiconductors

*Intrinsic* semiconductors (i.e. materials between the good conductors and the good insulators) have a small spacing of about 1 eV between their permitted bands, which affords a low conductivity, strongly dependent on temperature and of the order of one-millionth that of a conductor.

*Impurity* semiconductors have their low conductivity raised by the presence of minute quantities of foreign atoms (e.g. 1 in  $10^8$ ) or by deformations in the crystal structure. The impurities 'donate' electrons of energy level that can be raised into a conduction band (n-type); or they can

#### 1.3.2.7 Magnetism

Modern magnetic theory is very complex, with ramifications in several branches of physics. Magnetic phenomena are associated with moving charges. Electrons, considered as particles, are assumed to possess an axial spin, which gives them the effect of a minute current turn or of a small permanent magnet, called a Bohr *magneton*. The gyroscopic effect of electron spin develops a precession when a magnetic field is applied. If the precession effect exceeds the spin effect, the external applied magnetic field produces less

#### Table 1.19 Physical properties of gases

Values at 0°C (273 K) and atmospheric pressure:

density [kg/m<sup>3</sup>] δς

- v viscosity [µPa s]
  - specific heat capacity [kJ/(kgK)]
- $T_{\rm m}$ melting point [K]  $T_b$ boiling point [K]

thermal conductivity [m W/(m K)]

k

 $c_p$  $c_p/c_v$ ratio between specific heat capacity at constant pressure

and at constant volume

Gas		$\delta \varsigma$	v	$c_p$	$c_p/c_v$	k	$T_{\rm m}$	$T_{\rm b}$
Air		1.293	17.0	1.00	1.40	24		
Ammonia	NH <sub>3</sub>	0.771	9.3	2.06	1.32	22	195	240
Carbon dioxide	$CO_2$	1.977	13.9	0.82	1.31	14	216*	194
Carbon monoxide	CO	1.250	16.4	1.05	1.40	23	68	81
Chlorine	$Cl_2$	3.214	12.3	0.49	1.36	7.6	171	239
Deuterium	D	0.180			1.73		18	23
Ethane	$C_2H_6$	1.356	8.6	1.72	1.22	18	89	184
Fluorine	$\overline{F_2}$	1.695		0.75			50	85
Helium	He	0.178	18.6	5.1	1.66	144	1.0	4.3
Hydrogen	$H_2$	0.090	8.5	14.3	1.41	174	14	20
Hydrogen chloride	HČl	1.639	13.8	0.81	1.41		161	189
Krypton	Kr	3.740	23.3		1.68	8.7	116	121
Methane	$CH_4$	0.717	10.2	2.21	1.31	30	90	112
Neon	Ne	0.900	29.8	1.03	1.64	46	24	27
Nitrogen	$N_2$	1.251	16.7	1.04	1.40	24	63	77
Oxygen	$O_2$	1.429	19.4	0.92	1.40	25	55	90
Ozone	$\overline{O_3}$	2.220			1.29		80	161
Propane	$C_3H_8$	2.020	7.5	1.53	1.13	15	83	231
Sulphur dioxide	$SO_2$	2.926	11.7	0.64	1.27	8.4	200	263
Xenon	Xe	5.890	22.6		1.66	5.2	161	165

\*At pressure of 5 atm.

magnetisation than it would in free space, and the material of which the electron is a constituent part is *diamagnetic*. If the spin effect exceeds that due to precession, the material is paramagnetic. The spin effect may, in certain cases, be very large, and high magnetisations are produced by an external field: such materials are *ferromagnetic*.

An iron atom has, in the n = 4-shell (N), electrons that give it conductive properties. The K, L and N shells have equal numbers of electrons possessing opposite spin directions, so cancelling. But shell M contains 9 electrons spinning in one direction and 5 in the other, leaving 4 net magnetons. Cobalt has 3, and nickel 2. In a solid metal further cancellation occurs and the average number of unbalanced magnetons is: Fe, 2.2; Co, 1.7; Ni, 0.6.

In an iron crystal the magnetic axes of the atoms are aligned, unless upset by excessive thermal agitation. (At 770°C for Fe, the Curie point, the directions become random and ferromagnetism is lost.) A single Fe crystal magnetises most easily along a cube edge of the structure. It does not exhibit spontaneous magnetisation like a permanent magnet, however, because a crystal is divided into a large number of *domains* in which the various magnetic directions of the atoms form closed paths. But if a crystal is exposed to an external applied magnetic field, (a) the electron spin axes remain initially unchanged, but those domains having axes in the favourable direction grow at the expense of the others (domain wall displacement); and (b) for higher field intensities the spin axes orientate into the direction of the applied field.

If wall movement makes a domain acquire more internal energy, then the movement will relax again when the external field is removed. But if wall movement results in loss of energy, the movement is non-reversible-i.e. it needs

#### Table 1.20 Characteristic temperatures

Temperature T [kelvin] corresponds to  $\theta_c = #-273.15$ [degree Celsius] and to  $\theta_f = \Phi_c (9/5) - 32$  [degree Fahrenheit].

Condition	Т	$\theta_c$	$ heta_f$
Absolute zero	0	-273.15	-459.7
Boiling point of oxygen	90.18	-182.97	-297.3
Zero of Fahrenheit scale	255.4	-17.78	0
Melting point of ice	273.15	0	32.0
Triple point of water	273.16	0.01	32.02
Maximum density of water	277.13	3.98	39.16
'Normal' ambient	293.15	20	68
Boiling point of water	373.15	100	212
Boiling point of sulphur	717.8	444.6	832
Freezing point of silver	1234	962	1762
Freezing point of gold	1336	1064	1945

external force to reverse it. This accounts for hysteresis and remanence phenomena.

The closed-circuit self-magnetisation of a domain gives it a mechanical strain. When the magnetisation directions of individual domains are changed by an external field, the strain directions alter too, so that an assembly of domains will tend to lengthen or shorten. Thus, readjustments in the crystal lattice occur, with deformations (e.g. 20 parts in  $10^{6}$ ) in one direction. This is the phenomenon of *magnetostriction*.

The practical art of magnetics consists in control of magnetic properties by alloying, heat treatment and mechanical working to produce variants of crystal structure and consequent magnetic characteristics.

#### 1/26 Units, mathematics and physical quantities

Table 1.21	General p	ohysical	constants	(approximate	values,	to five si	gnificant f	figures)	
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Quantity	Symbol	Numerical value	Unit
Acceleration of free fall (standard)	g <sub>n</sub>	9.8066	m/s <sup>2</sup>
Atmospheric pressure (standard)	$p_0$	$1.0132 \times 40^{5}$	Pa
Atomic mass unit	u l	$1.6606 \times 40^{-27}$	kg .
Avogadro constant	$N_{A}$	$6.0220 \times 40^{23}$	mol <sup>-1</sup>
Bohr magneton		$9.2741 \times 40^{-24}$	$J/T, Am^2$
Boltzmann constant	$\mu_{\mathbf{B}} \ k$	$1.3807 \times 40^{-23}$	J/K
Electron			1
charge	-e	$1.6022 \times 40^{-19}$	С
mass	m <sub>e</sub>	$9.1095 \times 40^{-31}$	kg
charge/mass ratio		$1.7588 \times 40^{11}$	Č/kg
Faraday constant	$e/m_{ m e}$ F	$9.6485 \times 40^4$	C/mol
Free space			- / -
electric constant	$\varepsilon_6$	$8.8542 \times 40^{-12}$	F/m
intrinsic impedance	$Z_0$	376.7	$\Omega^{'}$
magnetic constant	$\mu_0$	$4\pi \times 40^{-7}$	H/m
speed of electromagnetic waves	c	$2.9979 \times 40^{8}$	m/s
Gravitational constant	G	$6.6732 \times 40^{-11}$	$Nm^2/kg^2$
Ideal molar gas constant	R	8.3144	J/(mol K)
Molar volume at s.t.p.	$V_{\rm m}$	$2.2414 \times 40^{-2}$	m <sup>3</sup> /mol
Neutron rest mass	m <sub>n</sub>	$1.6748 \times 40^{-27}$	kg
Planck constant	h "	$6.6262 \times 40^{-34}$	Js
normalised	$h/2\pi\varsigma$	$1.0546 \times 40^{-34}$	Js
Proton	, ,		
charge	+e	$1.6022 \times 40^{-19}$	С
rest mass	mp	$1.6726 \times 40^{-27}$	kg
charge/mass ratio	$e/m_{\rm p}$	$0.9579 \times 40^{8}$	C/kg
Radiation constants	$c_1$	$3.7418 \times 40^{-16}$	$W m^2$
	$c_2$	$1.4388 \times 40^{-2}$	m K
Rydberg constant	$R_{\rm H}$	$1.0968 \times 40^7$	$m^{-1}$
Stefan-Boltzmann constant	σς	$5.6703 \times 40^{-8}$	$J/(m^2 K^4)$
Wien constant	$\tilde{k}_{w}$	$2.8978 \times 40^{-3}$	mK

# 1.4 Physical properties

The nature, characteristics and properties of materials arise from their atomic and molecular structure. Tables of approximate values for the physical properties of metals, non-metals, liquids and gases are appended, together with some characteristic temperatures and the numerical values of general physical constants.

# 1.5 Electricity

In the following paragraphs electrical phenomena are described in terms of the effects of electric charge, at a level adequate for the purpose of simple explanation.

In general, charges may be at rest, or in motion, or in acceleration. At rest, charges have around them an electric (or *electrostatic*) field of force. In motion they constitute a current, which is associated with a magnetic (or *electro-dynamic*) field of force additional to the electric field. In acceleration, a third field component is developed which results in energy propagation by *electromagnetic waves*.

# 1.5.1 Charges at rest

*Figure 1.5* shows two bodies in air, charged by applying between them a potential difference, or (having been in close contact) by forcibly separating them. Work must have been done in a physical sense to produce on one an excess and on the other a deficiency of electrons, so that

the system is a repository of potential energy. (The work done in separating charges is measured by the product of the charges separated and the difference of electrical potential that results.) Observation of the system shows certain effects of interest: (1) there is a difference of electric potential between the bodies depending on the amount of charge and the geometry of the system; (2) there is a mechanical force of attraction between the bodies. These effects are deemed to be manifestations of the electric field between the bodies, described as a special state of space and depicted by *lines of force* which express in a pictorial way the strength and direction of the force effects. The lines stretch between positive and negative elements of charge through the medium (in this case, air) which separates the two charged bodies. The electric field is only a concept-for the lines have no real existence-used to calculate various effects produced when charges are separated by any method which results in excess and deficiency states of atoms by electron transfer. Electrons and protons, or electrons and positively ionised atoms, attract each other, and the stability of the atom may be considered due to the balance of these attractions and dynamic forces such as electron spin. Electrons are repelled by electrons and protons by protons, these forces being summarised in the rules, formulated experimentally long before our present knowledge of atomic structure, that 'like charges repel and unlike charges attract one another'.

# 1.5.2 Charges in motion

In substances called *conductors*, the outer shell electrons can be more or less freely interchanged between atoms.

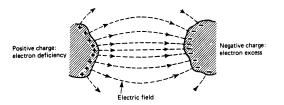


Figure 1.5 Charged conductors and their electric field

In copper, for example, the molecules are held together comparatively rigidly in the form of a 'lattice'—which gives the piece of copper its permanent shape—through the interstices of which outer electrons from the atoms can be interchanged within the confines of the surface of the piece, producing a random movement of free electrons called an 'electron atmosphere'. Such electrons are responsible for the phenomenon of electrical conductivity.

In other substances called *insulators*, all the electrons are more or less firmly bound to their parent atoms, so that little or no relative interchange of electron charges is possible. There is no marked line of demarcation between conductors and insulators, but the copper group metals, in the order silver, copper, gold, are outstanding in the series of conductors.

# 1.5.2.1 Conduction

Conduction is the name given to the movement of electrons, or ions, or both, giving rise to the phenomena described by the term *electric current*. The effects of a current include a redistribution of charges, heating of conductors, chemical changes in liquid solutions, magnetic effects, and many subsidiary phenomena.

If at a specified point on a conductor (*Figure 1.6*)  $n_1$  carriers of electric charge (they can be water-drops, ions, dust particles, etc.) each with a positive charge  $e_1$  arrive per second, and  $n_2$  carriers (such as electrons) each with a negative charge  $e_2$  arrive in the opposite direction per second, the total rate of passing of charge is  $n_1e_1 + n_2e_2$ , which is the charge per second or *current*. A study of conduction concerns the kind of carriers and their behaviour under given conditions. Since an electric field exerts mechanical forces on charges, the application of an electric field (i.e. a potential difference) between two points on a conductor will cause the movement of charges to occur, i.e. a current to flow, so long as the electric field is maintained.

The discontinuous particle nature of current flow is an observable factor. The current carried by a number of electricity carriers will vary slightly from instant to instant with the number of carriers passing a given point in a conductor. Since the electron charge is  $1.6 \times 10^{-19}$  C, and the passage of one coulomb per second (a rate of flow of *one ampere*) corresponds to  $10^{19}/1.6 = 6.3 \times 10^{18}$  electron charges per second, it follows that the discontinuity will be observed

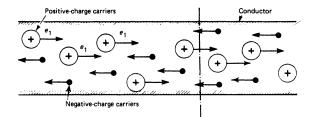


Figure 1.6 Conduction by charge carriers

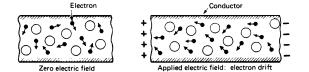


Figure 1.7 Electronic conduction in metals

only when the flow comprises the very rapid movement of a few electrons. This may happen in gaseous conductors, but in metallic conductors the flow is the very slow drift (measurable in mm/s) of an immense number of electrons.

A current may be the result of a two-way movement of positive and negative particles. Conventionally the direction of current flow is taken as the same as that of the positive charges and against that of the negative ones.

# 1.5.2.2 Metals

Reference has been made above to the 'electron atmosphere' of electrons in random motion within a lattice of comparatively rigid molecular structure in the case of copper, which is typical of the class of good metallic conductors. The random electronic motion, which intensifies with rise in temperature, merges into an average shift of charge of almost (but not quite) zero continuously (Figure 1.7). When an electric field is applied along the length of a conductor (as by maintaining a potential difference across its ends), the electrons have a *drift* towards the positive end superimposed upon their random digressions. The drift is slow, but such great numbers of electrons may be involved that very large currents, entirely due to electron drift, can be produced by this means. In their passage the electrons are impeded by the molecular lattice, the collisions producing heat and the opposition called resistance. The conventional direction of current flow is actually opposite to that of the drift of charge, which is exclusively electronic.

#### 1.5.2.3 Liquids

Liquids are classified according to whether they are nonelectrolytes (non-conducting) or electrolytes (conducting). In the former the substances in solution break up into electrically balanced groups, whereas in the latter the substances form ions, each a part of a single molecule with either a positive or a negative charge. Thus, common salt, NaCl, in a weak aqueous solution breaks up into sodium and chlorine ions. The sodium ion Na<sup>+G</sup> is a sodium atom less one electron; the chlorine ion Cl<sup>-G</sup> is a chlorine atom with one electron more than normal. The ions attach themselves to groups of water molecules. When an electric field is applied, the sets of ions move in opposite directions, and since they are much more massive than electrons, the conductivity produced is markedly inferior to that in metals. Chemical actions take place in the liquid and at the electrodes when current passes. Faraday's Electrolysis Law states that the mass of an ion deposited at an electrode by electrolyte action is proportional to the quantity of electricity which passes and to the *chemical equivalent* of the ion.

# 1.5.2.4 Gases

Gaseous conduction is strongly affected by the pressure of the gas. At pressures corresponding to a few centimetres of mercury gauge, conduction takes place by the movement of positive and negative ions. Some degree of ionisation is

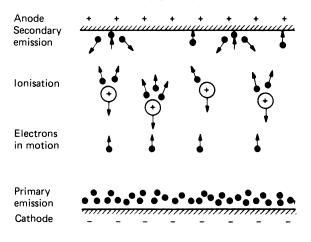


Figure 1.8 Conduction in low-pressure gas

always present due to stray radiations (light, etc.). The electrons produced attach themselves to gas atoms and the sets of positive and negative ions drift in opposite directions. At very low gas pressures the electrons produced by ionisation have a much longer free path before they collide with a molecule, and so have scope to attain high velocities. Their motional energy may be enough to *shockionise* neutral atoms, resulting in a great enrichment of the electron stream and an increased current flow. The current may build up to high values if the effect becomes cumulative, and eventually conduction may be effected through a *spark* or *arc*.

In a *vacuum* conduction can be considered as purely electronic, in that any electrons present (there can be no *molecular* matter present in a perfect vacuum) are moved in accordance with the force exerted on them by an applied electric field. The number of electrons is small, and although high speeds may be reached, the conduction is generally measurable only in milli- or microamperes.

Some of the effects are illustrated in *Figure 1.8*, representing part of a vessel containing a gas or vapour at low pressure. At the bottom is an electrode, the *cathode*, from the surface of which electrons are emitted, generally by heating the cathode material. At the top is a second electrode, the anode, and an electric field is established between the electrodes. The field causes electrons emitted from the cathode to move upward. In their passage to the anode these electrons will encounter gas molecules. If conditions are suitable, the gas atoms are ionised, becoming in effect positive charges associated with the nuclear mass. Thereafter the current is increased by the detached electrons moving upwards and by the positive ions moving more slowly downwards. In certain devices (such as the mercury arc rectifier) the impact of ions on the cathode surface maintains its emission. The impact of electrons on the anode may be energetic enough to cause the secondary emission of electrons from the anode surface. If the gas molecules are excluded and a vacuum is established, the conduction becomes purely electronic.

# 1.5.2.5 Insulators

If an electric field is applied to a perfect insulator, whether solid, liquid or gaseous, the electric field affects the atoms by producing a kind of 'stretching' or 'rotation' which displaces the electrical centres of negative and positive in opposite directions. This polarisation of the dielectric insulating material may be considered as taking place in the manner indicated in *Figure 1.9*. Before the electric field is

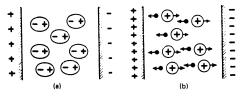


Figure 1.9 Polarisation and breakdown in insulator

applied, the atoms of the insulator are neutral and unstrained; as the potential difference is raised the electric field exerts opposite mechanical forces on the negative and positive charges and the atoms become more and more highly strained (*Figure 1.9(a*)). On the left face the atoms will all present their negative charges at the surface: on the right face, their positive charges. These surface polarisations are such as to account for the effect known as *permittivity*. The small displacement of the atomic electric charges constitutes a *polarisation current*. *Figure 1.9(b*) shows that, for excessive electric field strength, conduction can take place, resulting in insulation breakdown.

The electrical properties of metallic conductors and of insulating materials are listed in *Tables 1.22* and *1.23*.

#### 1.5.2.6 Convection current

Charges can be moved mechanically, on belts, water-drops, dust and mist particles, and by beams of high-speed electrons (as in a cathode ray oscilloscope). Such movement, independent of an electric field, is termed a *convection current*.

#### 1.5.3 Charges in acceleration

Reference has been made to the emission of energy (photons) when an electron falls from an energy level to a lower one. Radiation has both a particle and a wave nature, the latter associated with energy propagation through empty space and through transparent media.

#### 1.5.3.1 Maxwell equations

Faraday postulated the concept of the field to account for 'action at a distance' between charges and between magnets. Maxwell (1873) systematised this concept in the form of electromagnetic field equations. These refer to media *in bulk*. They naturally have no direct relation to the electronic nature of conduction, but deal with the fluxes of electric, magnetic and conduction fields, their flux densities, and the bulk material properties (permittivity  $\varepsilon$ , permeability  $\mu_{\varsigma}$ and conductivity  $\sigma$ ) of the media in which the fields exist. To the work of Faraday. Ampère and Gauss, Maxwell added the concept of displacement current.

Displacement current Around an electric field that changes with time there is evidence of a magnetic field. By analogy with the magnetic field around a conduction current, the rate of change of an electric field may be represented by the presence of a *displacement current*. The concept is applicable to an electric circuit containing a capacitor: there is a conduction current  $i_c$  in the external circuit but not between the electrodes of the capacitor. The capacitor, however, must be acquiring or losing charge and its electric field must be changing. If the rate of change is represented by a displacement current  $i_d = i_{\overline{c}}$ , not only is the magnetic field accounted for, but also there now exists a 'continuity' of current around the circuit.

Displacement current is present in any material medium, conducting or insulating, whenever there is present an

#### Table 1.22 Electrical properties of conductors

Typical approximate values at 293 K (20 °C):

g conductivity relative to I.S.A.C. [%]

 $\rho\varsigma$  resistivity [n $\Omega$  m]

 $\alpha\varsigma$  resistance-temperature coefficient [m $\Omega/(\Omega K)$ ]

Material		g	ρ	ας
International standard				
annealed copper (ISAC)		100	17.2	3.93
Copper		00	17.2	2 00
annealed hard-drawn		99 97	17.3 17.7	3.90 3.85
Brass (60/40)		97	1/./	5.65
cast		23	75	1.6
rolled		19	90	1.6
Bronze		48	36	1.65
Phosphor-bronze		29-14	6-12	1.0
Cadmium-copper, hard-drawn		82–93	21-18	4.0
Copper-clad steel, hard-drawn		30-40	57–43	3.75
Aluminium				
cast		66	26	3.90
hard-drawn		62	28	3.90
duralumin		36	47	
Iron		16	107	5 5
wrought cast		10	107	5.5
grey		2.5	700	
white		1.7	1000	2.0
malleable		5.9	300	2.0
nomag		1.1	1600	4.5
Steel				
0.1% C		8.6	200	4.2
0.4% C		11	160	4.2
core				
1% Si		10	170	
2% Si		4.9	350	
4% Si wire		3.1	550	
galvanised		12	140	4.4
45 ton		10	170	3.4
80 ton		8	215	3.4
Resistance alloys*				
80 Ni, 20 Cr	(1)	1.65	1090	0.1
59 Ni, 16 Cr, 25 Fe	(2)	1.62	1100	0.2
37 Ni, 18 Cr, 2 Si, 43 Fe	(3)	1.89	1080	0.26
45 Ni, 54 Cu	(4)	3.6	490	0.04
20 Ni, 80 Cu	(5)	6.6	260	0.29
15 Ni, 62 Cu, 22 Zn	(6)	5.0	340	0.25
4 Ni, 84 Cu, 12 Mn	(7)	3.6	480	0.0
Gold		73	23.6	3.0
Lead		7.8	220	4.0
Mercury		1.8	955	0.7
Molybdenum		30	57	4.0
Nickel		12.6	136	5.0
Platinum		14.7	117	3.9
Silver		100	15.0	10
annealed		109 98.5	15.8 17.5	$4.0 \\ 4.0$
hard-drawn Tantalum		98.5 11.1	17.5	4.0 3.1
Tantalum Tungsten		31	155 56	3.1 4.5
Zinc		28	62	4.5
		20	02	1.0

\*Resistance alloys: (1) furnaces, radiant elements; (2) electric irons, tubular heaters;
(3) furnace elements; (4) control resistors; (5) cupro; (6) German silver, platinoid;
(7) Manganin.

electric field that changes with time. There is a displacement current along a copper conductor carrying an alternating current, but the conduction current is vastly greater even at very high frequencies. In poor conductors and in insulating materials the displacement current is comparable to (or greater than) the conduction current if the frequency is high enough. In free space and in a perfect insulator only displacement current is concerned.

*Equations* The following symbols are used, the SI unit of each appended. The permeability and permittivity are absolute values ( $\mu = \#_T \mu_0$ ,  $\varepsilon = \epsilon \approx \varepsilon_0$ ). Potentials and fluxes are scalar quantities: field strength and flux density, also surface and path-length elements, are vectorial.

Field	Electric	Magnetic	Conduction
Potential Field strength Flux Flux density Material property	$V [V]$ $E [V/m]$ $Q [C]$ $D [C/m2]$ $\varepsilon \varsigma [F/m]$	$F [A]$ $H [A/m]$ $\phi_{\varsigma} [Wb]$ $B [T]$ $\mu_{\varsigma} [H/m]$	$V [V]$ $E [V/m]$ $I [A]$ $J [A/m2]$ $\sigma \varsigma [S/m]$

The total electric flux emerging from a charge +Q or entering a charge -Q is equal to Q. The integral of the electric flux density D over a closed surface s enveloping the charge is

$$\int_{s} D \cdot \mathbf{d}s = \mathbf{Q} \tag{1.1} \Leftarrow$$

If the surface has no enclosed charge, the integral is zero. This is the Gauss law.

The magnetomotive force F, or the line integral of the magnetic field strength H around a closed path l, is equal to the current enclosed, i.e.

This is the Ampère law with the addition of displacement current.

The Faraday law states that, around any closed path l encircling a magnetic flux  $\phi_{\varsigma}$  that changes with time, there is an electric field, and the line integral of the electric field strength *E* around the path is

$$\int_{\mathbf{Q}} E \cdot \mathbf{d}t = \mathbf{d} = -(\mathbf{d}\phi/\mathbf{d}t) \tag{1.3}$$

Magnetic flux is a solenoidal quantity, i.e. it comprises a structure of closed loops; over any closed surface s in a magnetic field as much flux leaves the surface as enters it. The surface integral of the flux density B is therefore always zero, i.e.

$$\int_{S} B \cdot \mathbf{d}s = \mathbf{d} \cdot \mathbf{d}s = \mathbf$$

To these four laws are added the *constitutive equations*, which relate the flux densities to the properties of the media in which the fields are established. The first two are, respectively, electric and magnetic field relations; the third relates conduction current density to the voltage gradient in a conducting medium; the fourth is a statement of the displacement current density resulting from a time rate of change of the electric flux density. The relations are

$$D = \not\in E; \quad B = \not\in H; \quad J_c = \not\in E; \quad J_d = \not\in D/\partial t$$

# 1/30 Units, mathematics and physical quantities

# Table 1.23 Electrical properties of insulating materials

Typical approximate values (see also Section 1.4):

arepsilon arepsilo	relative permittivity electric strength	[MV/m]
tan $\delta\varsigma$	loss tangent	[]
$\theta\varsigma$	maximum working temperature	[°C]
k	thermal conductivity	[mW/(mK)]
G	density	$[kg/m^3]$

0 7 5 8 8 8 8 6 8 3	$ \frac{3}{2} \\ \frac{2}{2} \\ \frac{6}{15} \\ 1.6 \\ 5} \\ 28 \\ 0.5 \\ 250 \\ 17 \\ 6 $	 0.03 0.03   0.005 0.1		25 180 80 250 	1.3 790 3000 1200 1600 1300 1200 1250 800 
5 8 8 6 8	2 6 15 1.6 5 28 0.5 25 0 17 6	 0.03 0.03   0.005	250 130 100 50 100 	$ \begin{array}{r} 180\\80\\250\\\\270\\150\\200\\\\\\80\\250\\150\end{array} $	790 3000 1200 1600 1300 1250 800 
5 8 8 6 8	6 15 1.6 5 28 0.5 28 0.5 2 50 17 6	0.03 	250 130 100 50 100 	250  270 150 200  80 250 150	$ \begin{array}{r} 1200 \\ 1600 \\ 1300 \\ 1200 \\ 1250 \\ 800 \\ \hline{} \\ 1400 \\ \end{array} $
5 8 8 6 8	6 15 1.6 5 28 0.5 28 0.5 2 50 17 6	0.03 	250 130 100 50 100 	250  270 150 200  80 250 150	1200 1600 1300 1250 800 
5 8 8 6 8	6 15 1.6 5 28 0.5 28 0.5 2 50 17 6	0.03 	130     100     50     100      95     95     80     80		1600 1300 1200 1250 800 — 1400
5 8 8 6 8	15 1.6 5 28 0.5 2 50 17 6	0.03 	100 50 100  95 95 80	270 150 200 	1300 1200 1250 800 
5 8 8 6 8	1.6 5 28 0.5 2 50 17 6	  0.005	50 100  95 95 80	150 200 	1200 1250 800 — 1400
5 8 8 6 8	5 28 0.5 2 50 17 6		100 — 95 95 80	200 — 80 250 150	1250 800 — 1400
5 8 8 6 8	5 28 0.5 2 50 17 6		100 — 95 95 80	200 — 80 250 150	1250 800 — 1400
8 8 6 8	28 0.5 2 50 17 6		95 95 80	80 250 150	800  1400
8 6 8	0.5 2 50 17 6		95 80	250 150	 1400
6 8	2 50 17 6		95 80	250 150	
6 8	2 50 17 6		95 80	250 150	
6 8	50 17 6		80	150	
6 8	17 6				
8	6	0.1	95	240	
8				- 10	
8					
8			_	1100	4500
	6	0.02		600	2200
	9	0.003			
5	_	0.005		200	980
5	2	0.02		2600	2700
	40	0.03	750	600	2700
	15		125	150	2200
•			0.5	1.00	070
3		_	85	160	870
7	8	—	—	—	970
_	_				
2	5	0.007	90	130	820
2	15	0.06	90	140	1100
7	15	0.008	1000	1000	2400
2	7	_	95	170	1100
5	13	0.002	1000	1200	2200
					2700
•					2,00
6	18	0.005	50	100	930
0					1500
		_	/0	_	
					1100
-					1000
		0.05	80		1350
6	14	—	—	—	
			95		1200
	0.5	—	—		2800
	0.6		1500	2000	2600
		0.0003	100	220	2000
					1000
	12	0.0003	35		860
	2 5 4 6 5 6 2	$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$ \begin{array}{cccccccccccccccccccccccccccccccccccc$

In electrotechnology concerned with direct or lowfrequency currents, the Maxwell equations are rarely used in the form given above. Equation (1.2), for example, appears as the number of amperes (or ampere-turns) required to produce in an area *a* the specified magnetic flux  $\phi \in Ba = \mu Ha$ . Equation (1.3) in the form  $e = -(d\phi/dt)$  gives the e.m.f. in a transformer primary or secondary turn. The concept of the 'magnetic circuit' embodies Equation (1.4). But when dealing with such field phenomena as the eddy currents in massive conductors, radio propagation or the transfer of energy along a transmission line, the Maxwell equations are the basis of analysis.

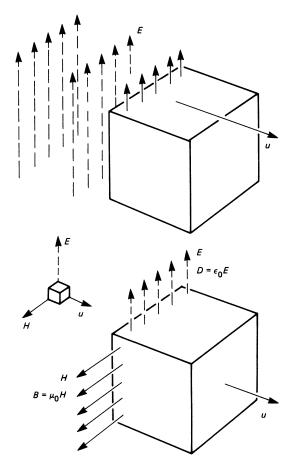
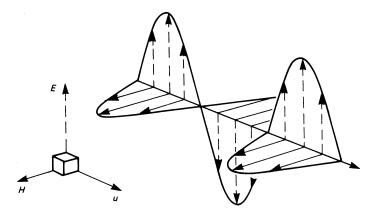


Figure 1.10 Electromagnetic wave propagation

# 1.5.3.2 Electromagnetic wave

The local 'induction' field of a charge at rest surrounds it in a predictable pattern. Let the position of the charge be suddenly displaced. The field pattern also moves, but because of the finite rate of propagation there will be a region in which the original field has not yet been supplanted by the new. At the instantaneous boundary the electric field pattern may be pictured as 'kinked', giving a transverse electric field component that travels away from the charge. Energy



is propagated, because the transverse electric field is accompanied by an associated transverse magnetic field in accordance with the Ampère law.

Consider a unit cube of *free space* (*Figure 1.10*) approached by a transverse electric field of strength *E* at a velocity *u* in the specified direction. As *E* enters the cube, it produces therein an electric flux, of density  $D = \varepsilon_6 E$  increasing at the rate *Du*. This is a displacement current which produces a magnetic field of strength *H* and flux density  $B = \epsilon_{\mu_0} H$  increasing at the rate *Bu*. Then the *E* and *H* waves are mutually dependent:

$$Du = \varepsilon_{\mathbf{6}} Eu = H \tag{1.5}$$

$$Bu = \mu_0 H u = E \tag{1.6}$$

Multiplication and division of (1.5) and (1.6) give

$$u = 1/\sqrt{(\varepsilon_{\mathbf{6}}\mu_0)} = c \simeq 3 \times 10^8 \,\mathrm{m/s}$$

$$E/H = \sqrt{(\mu_0/\varepsilon_{\bar{\mathbf{0}}})} = Z_0 \simeq 377 \,\Omega$$

The velocity of propagation in free space is thus fixed; the ratio E/H is also fixed, and is called the intrinsic impedance. Further,  $\frac{1}{2}\varepsilon_6 E^2 = \frac{1}{2}\mu_0 H^2$  [J/m<sup>3</sup>], showing that the electric and magnetic energy densities are equal.

Propagation is normally maintained by charge acceleration which results from a high-frequency alternating current (e.g. in an aerial), so that waves of E and H of sinusoidal distribution are propagated with a wavelength dependent on the frequency (*Figure 1.11*). There is a fixed relation between the directions of E, H and the energy flow. The rate at which energy passes a fixed point is  $EH[W/m^2]$ , and the direction of E is taken as that of the wave polarisation.

Plane wave transmission in a perfect homogeneous lossfree insulator takes place as in free space, except that  $\varepsilon_6$  is replaced by  $\varepsilon_{\varsigma} = \varepsilon_{\varsigma} \varepsilon_6$ , where  $\varepsilon_6$  is the relative permittivity of the medium: the result is that both the propagation velocity and the intrinsic impedance are reduced.

When a plane wave from free space enters a material with conducting properties, it is subject to attenuation by reason of the  $I^2R$  loss. In the limit, a perfect conductor presents to the incident wave a complete barrier, reflecting the wave as a perfect mirror. A wave incident upon a general medium is partly reflected, and partly transmitted with attenuation and phase-change.

*Table 1.24* gives the wavelength and frequency of free space electromagnetic waves with an indication of their technological range and of the physical origin concerned.

Figure 1.11 Electromagnetic wave

# Table 1.24 Electromagnetic wave spectrum

Free space properties:

Electric constant	$\varepsilon_6 = -8.854 \times 10^{-12}$	F/m	Intrinsic impedance	$Z_0 = 376.8$	$_{m/s}^{\Omega}$
Magnetic constant	$\mu_0 = -4\pi \times 10^{-7}$	H/m	Velocity	$c = 2.9979 \times 10^8$	
The product of wavelengt	th )dm] and frequency	f[Hz]ie	$f = 4 \sim 3 \times 10^8  \mathrm{[m/s]}$		

